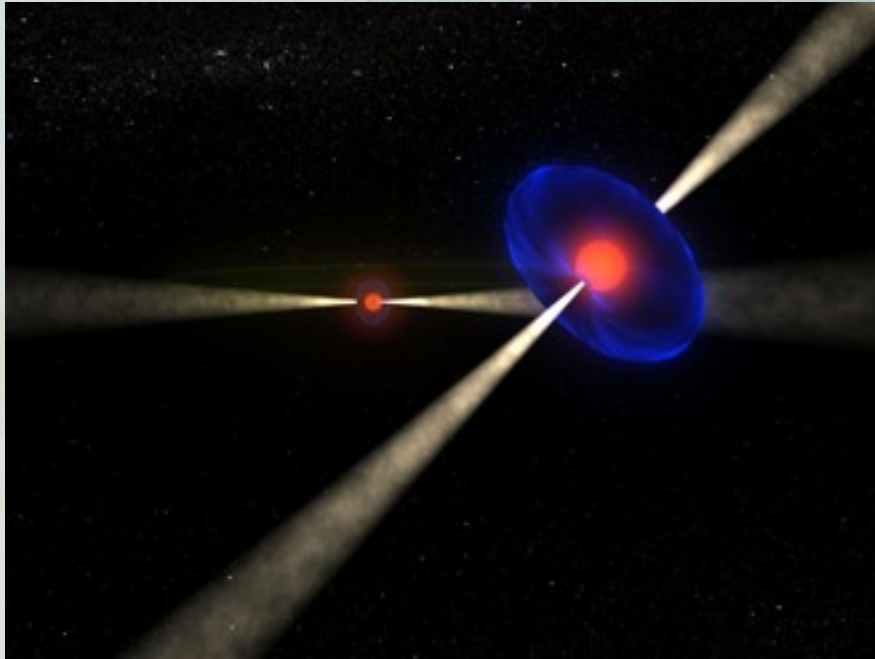


# Binary Pulsars as Tools to Study Gravity



© Daniel Cantin / McGill University

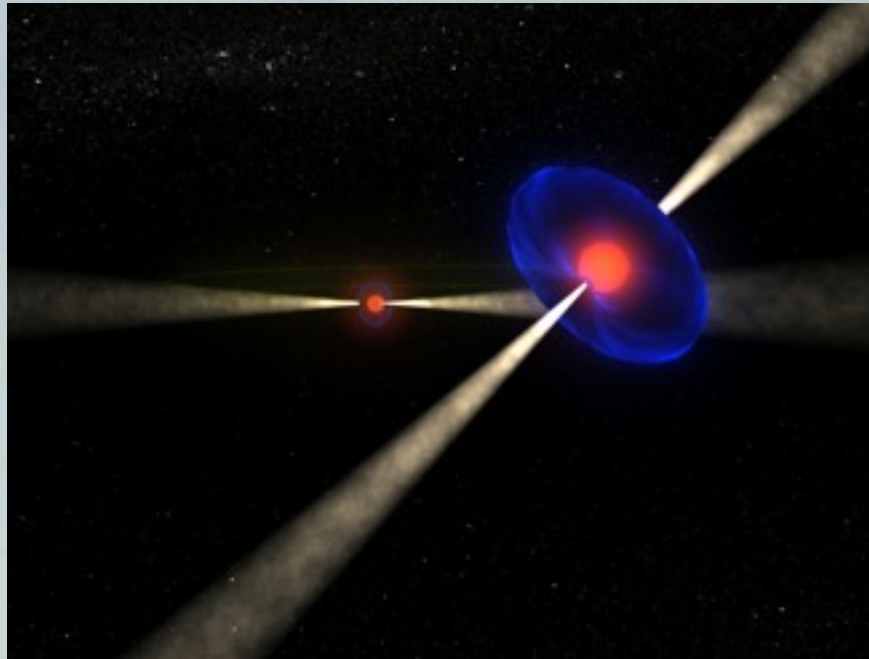
**René Breton**  
**University of Southampton**

GWPAW 2012  
Hannover, 7 June 2012

# By the End You Should Remember...

## Binary pulsars

- Excellent tools to **test gravity**
- Probe a **different** gravitational field **regime**
- **System diversity** = complementarity



© Daniel Cantin / McGill University

# Pulsars 101

# The Binary Pulsar Population

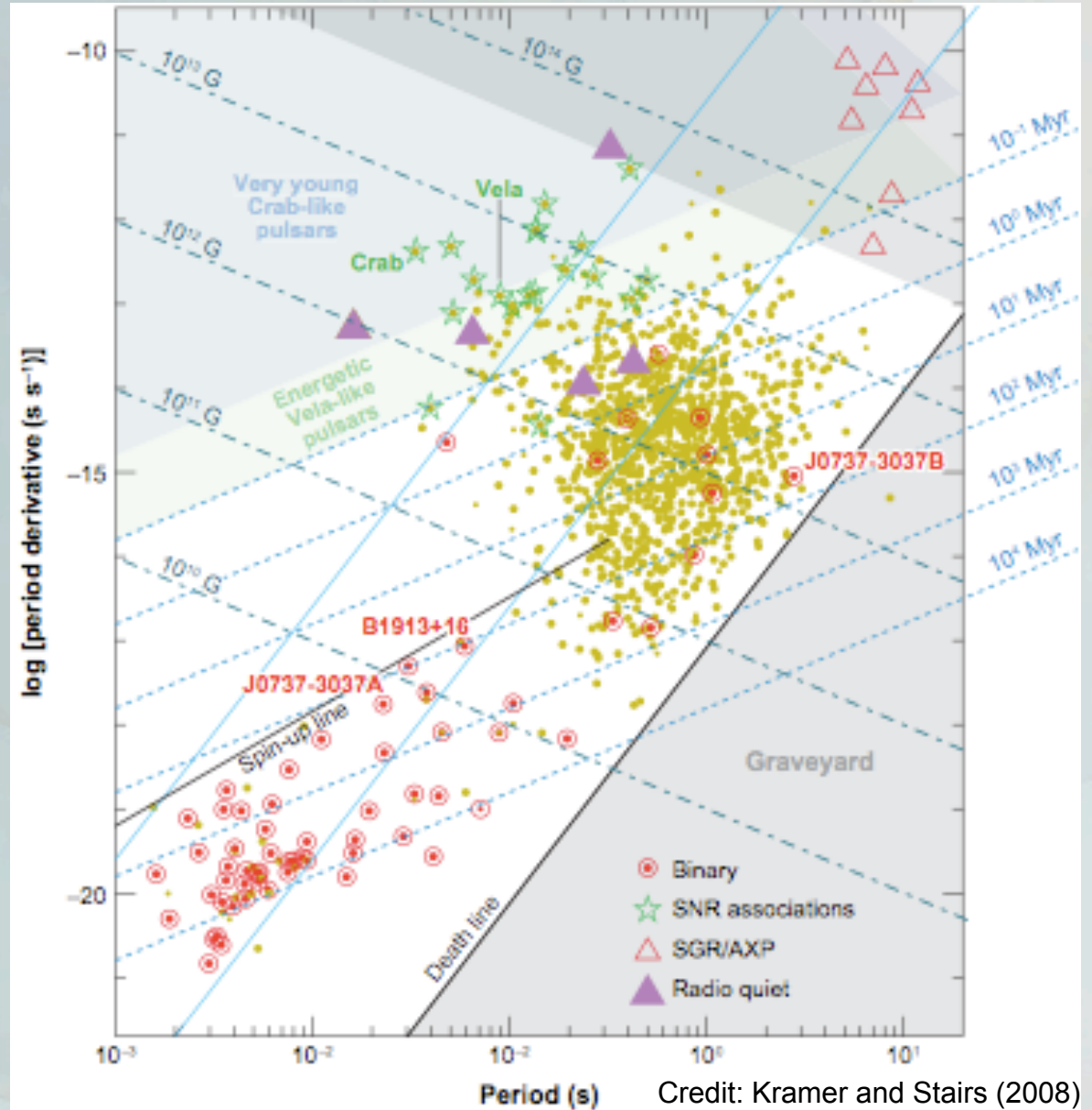
Binary pulsars are “special”

According to the ATNF catalogue:

- 2008 pulsars
- 186 binary pulsars

Binaries:

- **Short** spin periods
- **Old**
- **Low** magnetic fields



# The Binary Pulsar Population

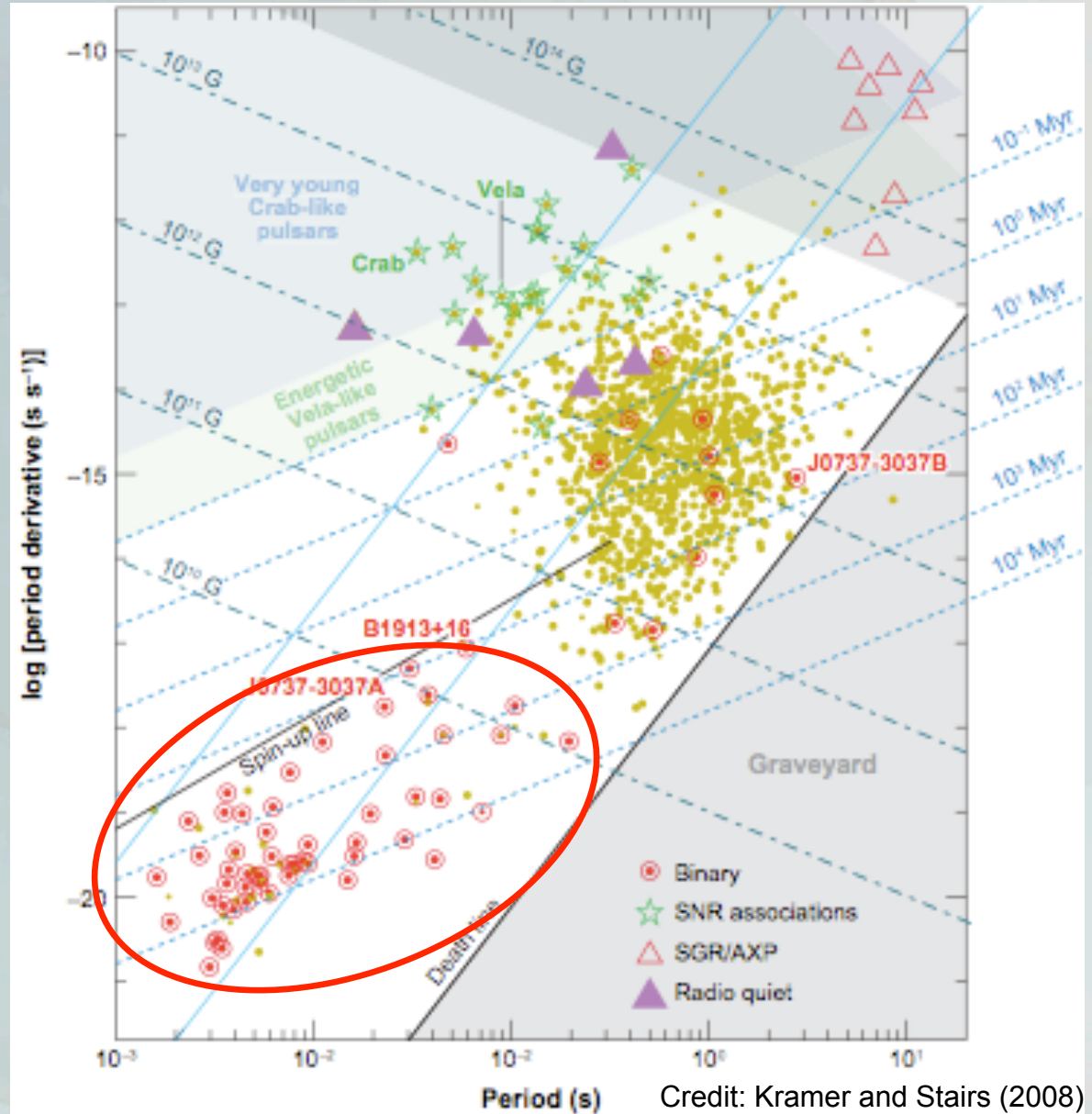
Binary pulsars are “special”

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# Binary Pulsar Timing and Gravity

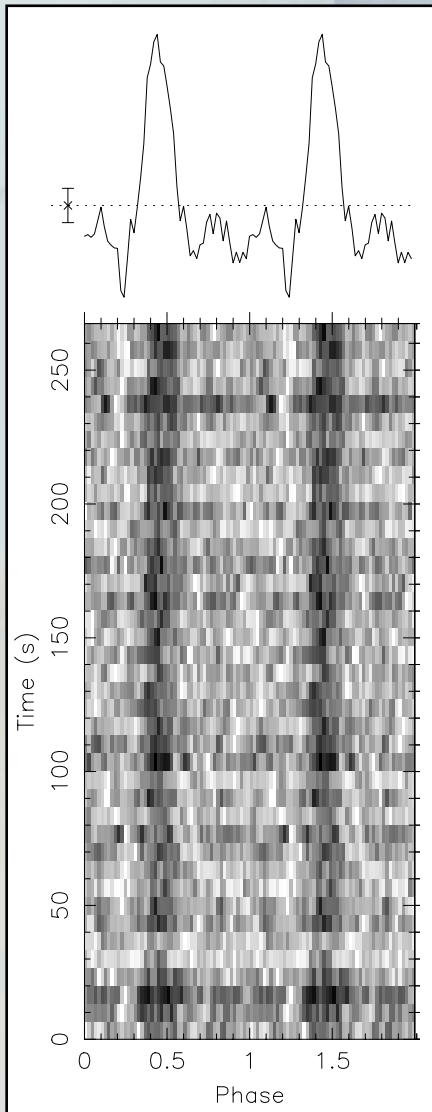
Almost textbook examples

*Compact objects (point mass)*

*Precise timing (frequency standard)*

Typical TOAs precision is  $\sim 0.1\% P_{\text{spin}}$

$$\phi = \phi_0 + \nu(t - t_0) + \frac{1}{2}\dot{\nu}(t - t_0)^2 + \dots$$



# Binary Pulsar Dynamics & Relativistic Effects

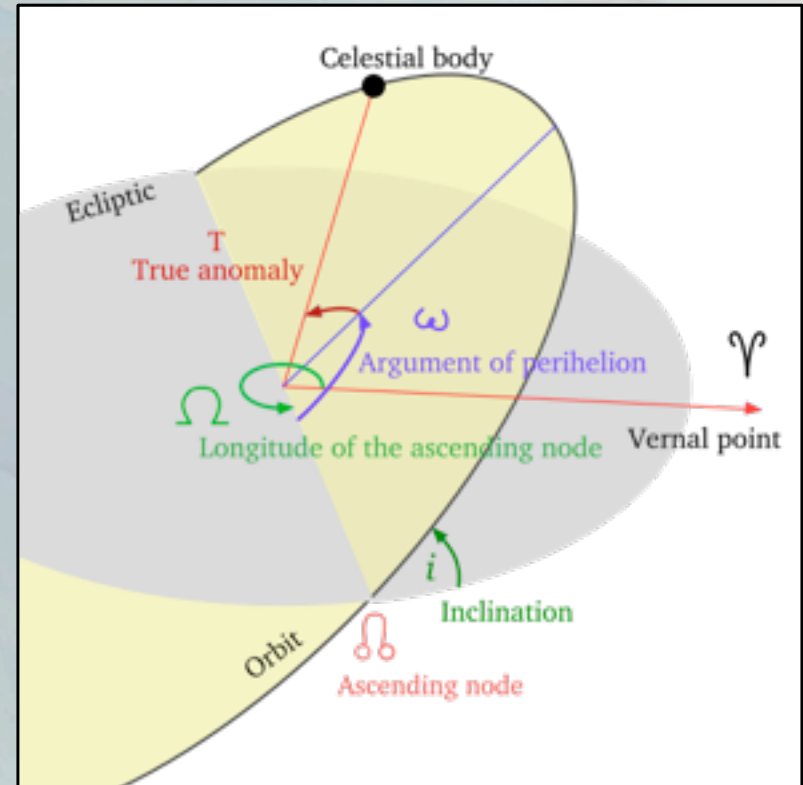
# Binary Pulsar Timing

## Newtonian orbits

Timing binary pulsars allows one to determine 5 Keplerian parameters:

- Orbital period  $P_b$
- Eccentricity  $e$
- Longitude of periastron  $\omega$
- Projected semimajor axis  $x_1 \equiv a_1/c \sin i$
- Epoch of periastron  $T_0$

We measure radial Doppler shifts only...  
=> The **orbital inclination** angle is **unknown**.



Arpad Horvath (Wikipedia)



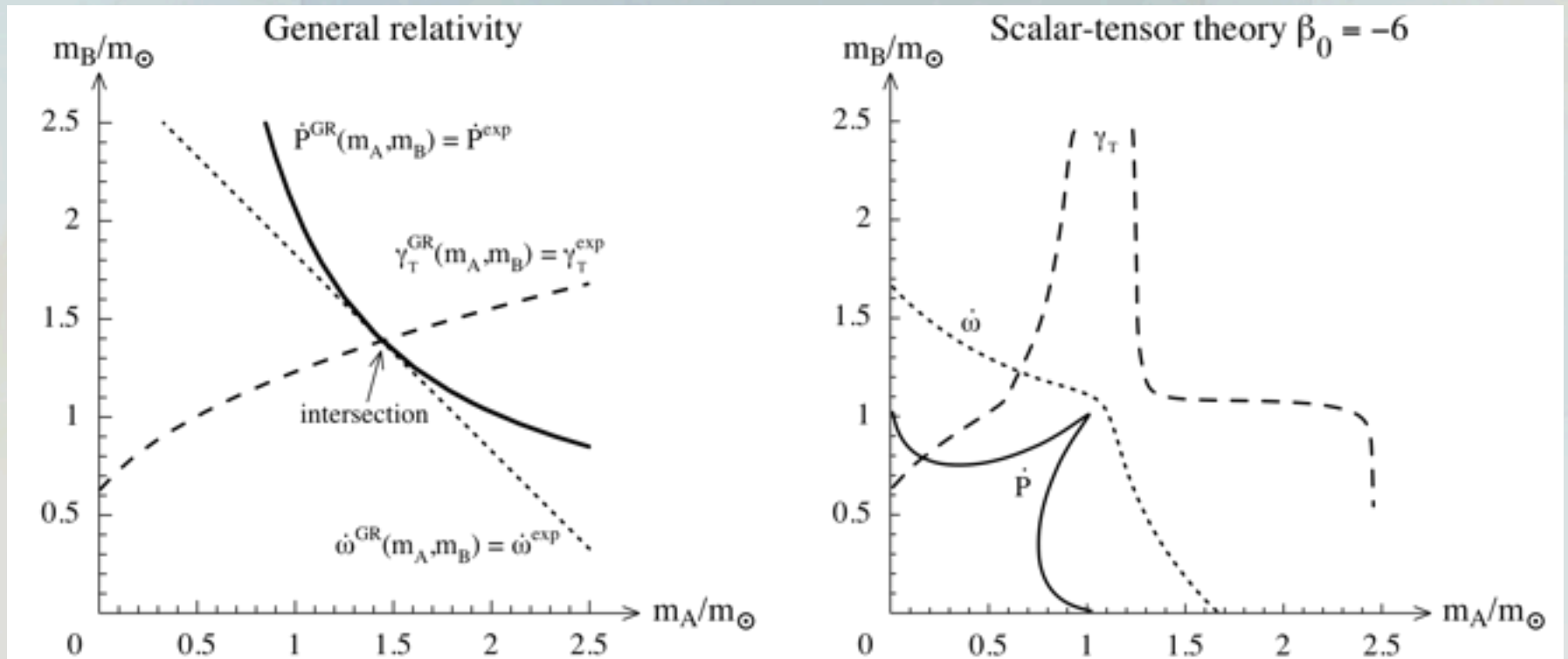
# Binary Pulsar Timing

## Relativistic orbits

**Post-Keplerian** (PK) parameters: dynamics beyond Newton.

- PKs are **phenomenological corrections** (theory-independent).
- PKs are functions of Keplerian parameters,  $M_1$  and  $M_2$ .

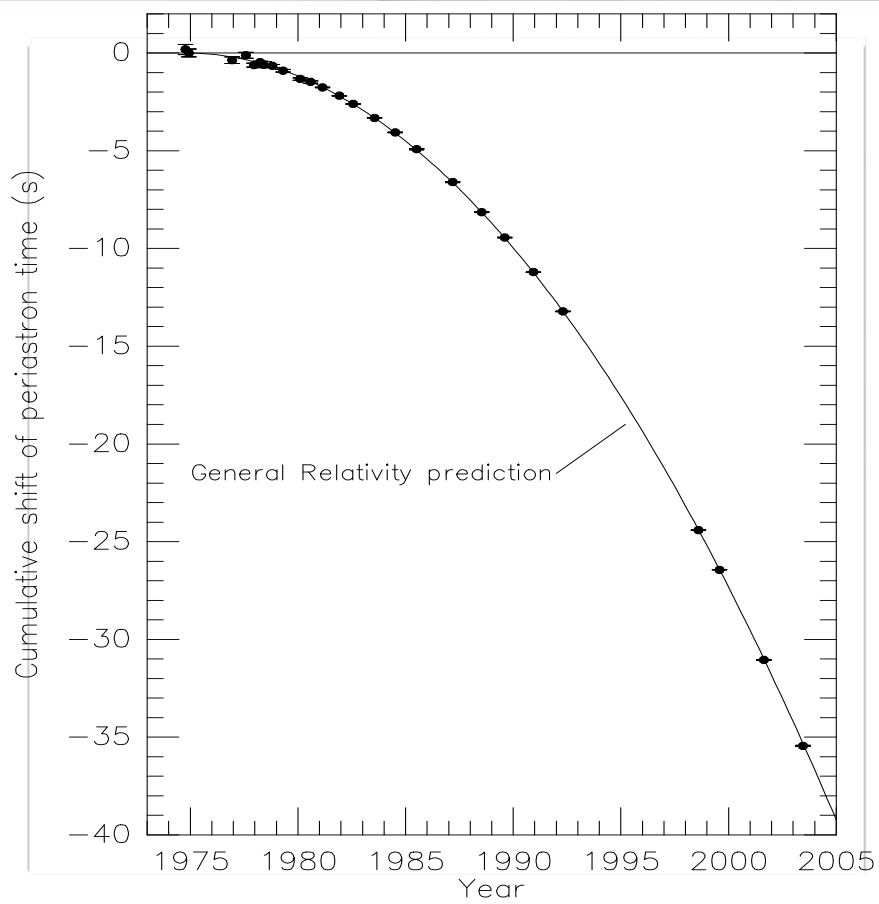
More than 2 PKs = test of a given relativistic theory of gravity.



Esposito-Farese (2004)

# Binary Pulsar Timing

## Relativistic orbits



Weisberg & Taylor (2005)

- Testing gravity relies on **orbital dynamics**
- Ideally: pulsar + neutron star  
eccentric + compact orbit
- About 10 known “relativistic” pulsar binaries

### PSR B1913+16

- First binary pulsar discovered (Hulse & Taylor 1975)
- First indirect evidence of gravitational wave emission (Taylor et al. 1979)

# The Double Pulsar

## A unique system

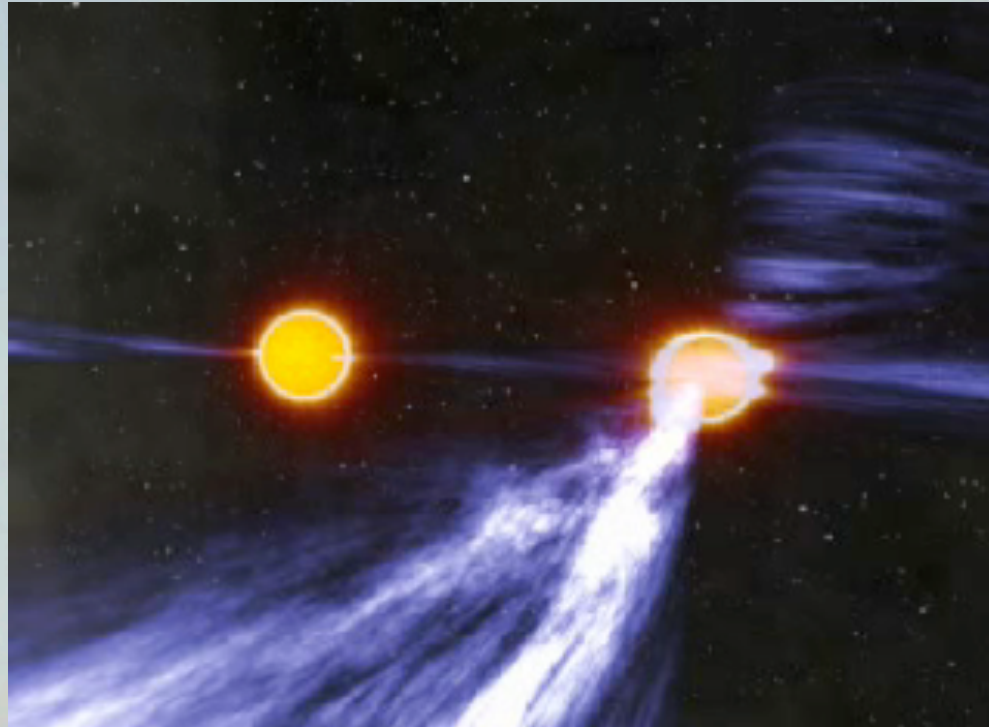
The **Double Pulsar** (PSR J0737-3039A/B), is the **only known pulsar-pulsar** system.

Pulsar A: 23 ms (Burgay et al., 2003)

Pulsar B: 2.8 s (Lyne et al., 2004)

Eccentricity: 0.08

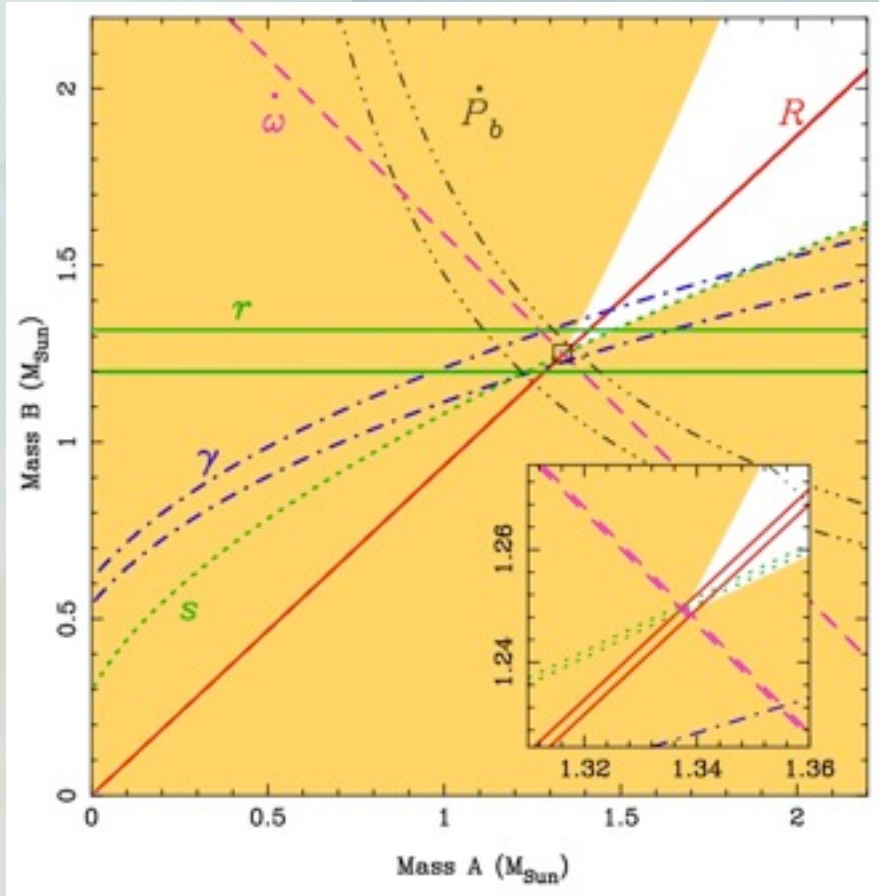
Orbital period: 2.4 hrs !!! ( $v_{\text{orbital}} \sim 0.001 c$ )



© John Rowe Animation  
Australia Telescope National Facility, CSIRO

# The Double Pulsar

## Testing gravity



Kramer et al., 2006

$$\dot{\omega} = 3T_{\odot}^{2/3} \left( \frac{P_b}{2\pi} \right)^{-5/3} \frac{1}{1-e^2} (M_A + M_B)^{2/3},$$

$$\gamma = T_{\odot}^{2/3} \left( \frac{P_b}{2\pi} \right)^{1/3} e \frac{M_B(M_A + 2M_B)}{(M_A + M_B)^{4/3}},$$

$$\dot{P}_b = -\frac{192\pi}{5} T_{\odot}^{5/3} \left( \frac{P_b}{2\pi} \right)^{-5/3} \frac{(1 + \frac{71}{24}e^2 + \frac{37}{96}e^4)}{(1-e^2)^{7/2}} \frac{M_A M_B}{(M_A + M_B)^{1/3}},$$

$$r = T_{\odot} M_B,$$

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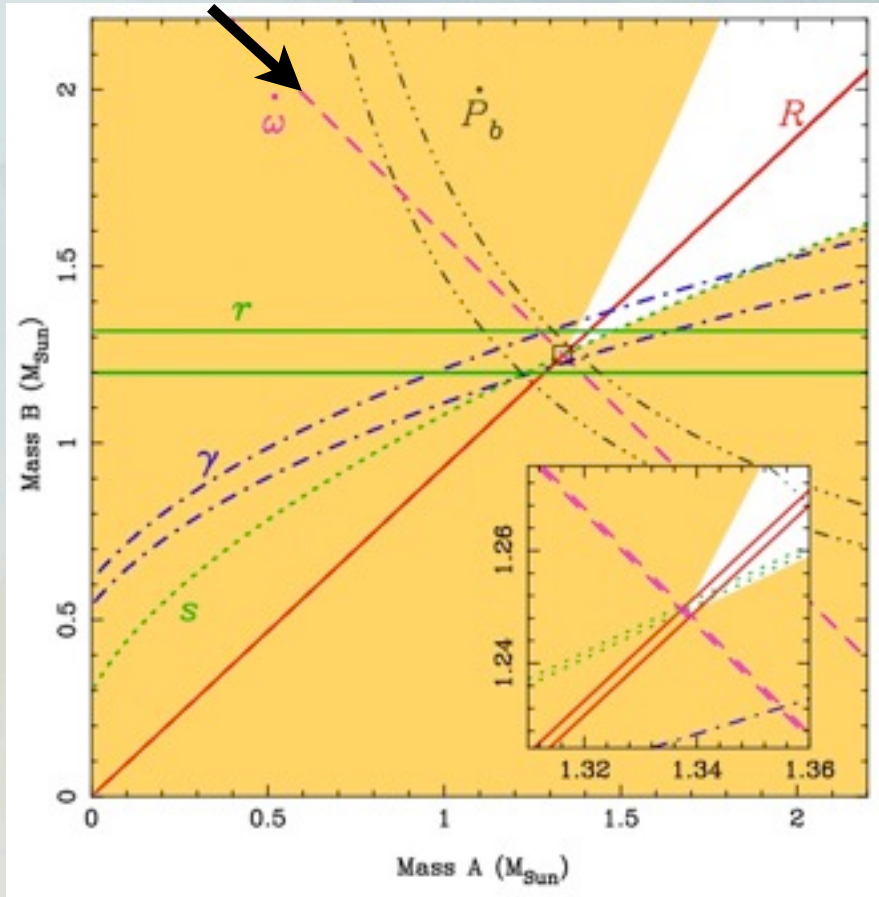
## Periastron advance

Rate of precession of the periastron:  $17^{\circ}\text{yr}^{-1}$  !

It takes 21 years to complete a cycle of precession.

# The Double Pulsar

## Testing gravity



Kramer et al., 2006

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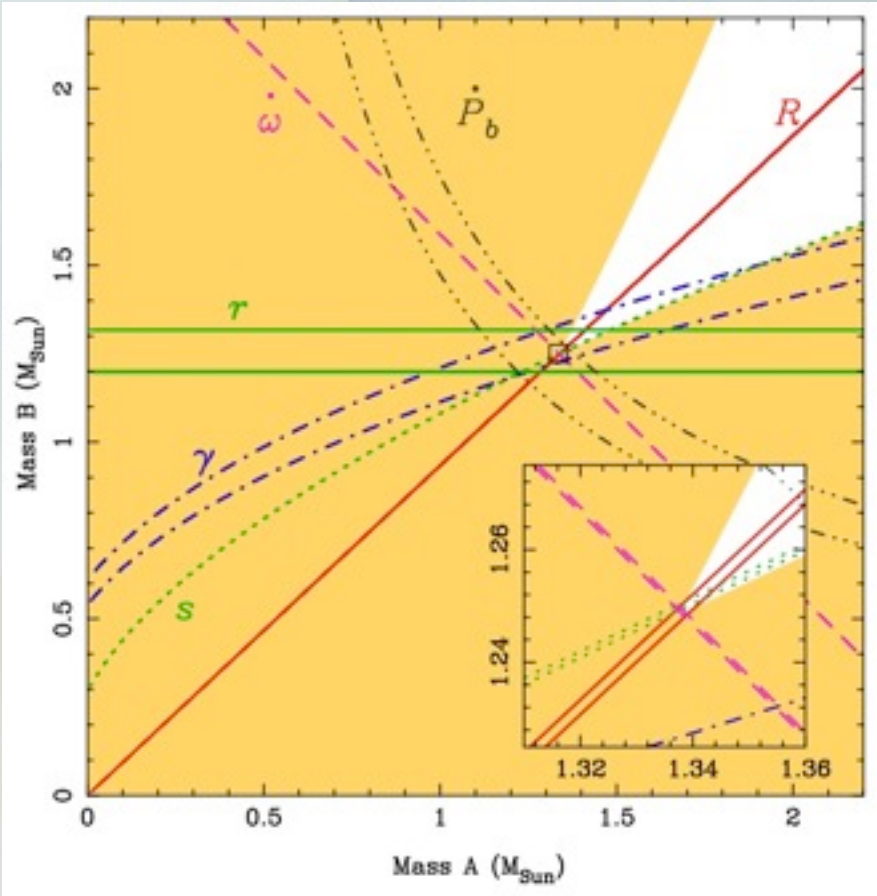
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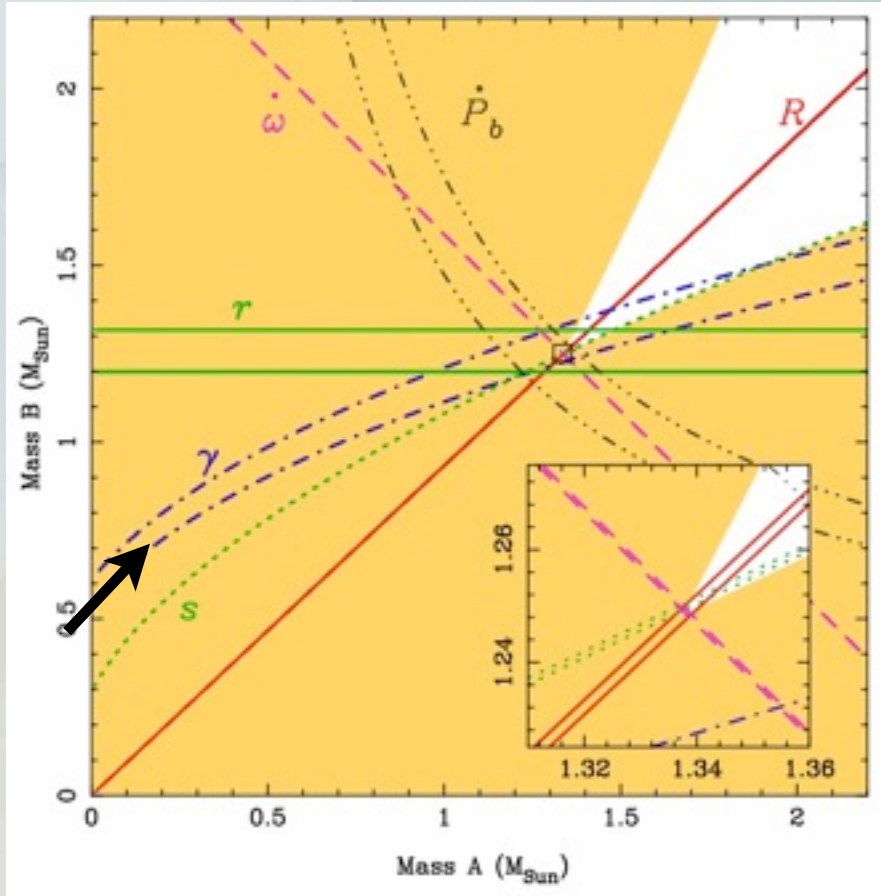
## Gravitational redshift and time dilation

Needs to climb to potential well

Clocks tick at variable rate in gravitational potentials

# The Double Pulsar

## Testing gravity



Kramer et al., 2006

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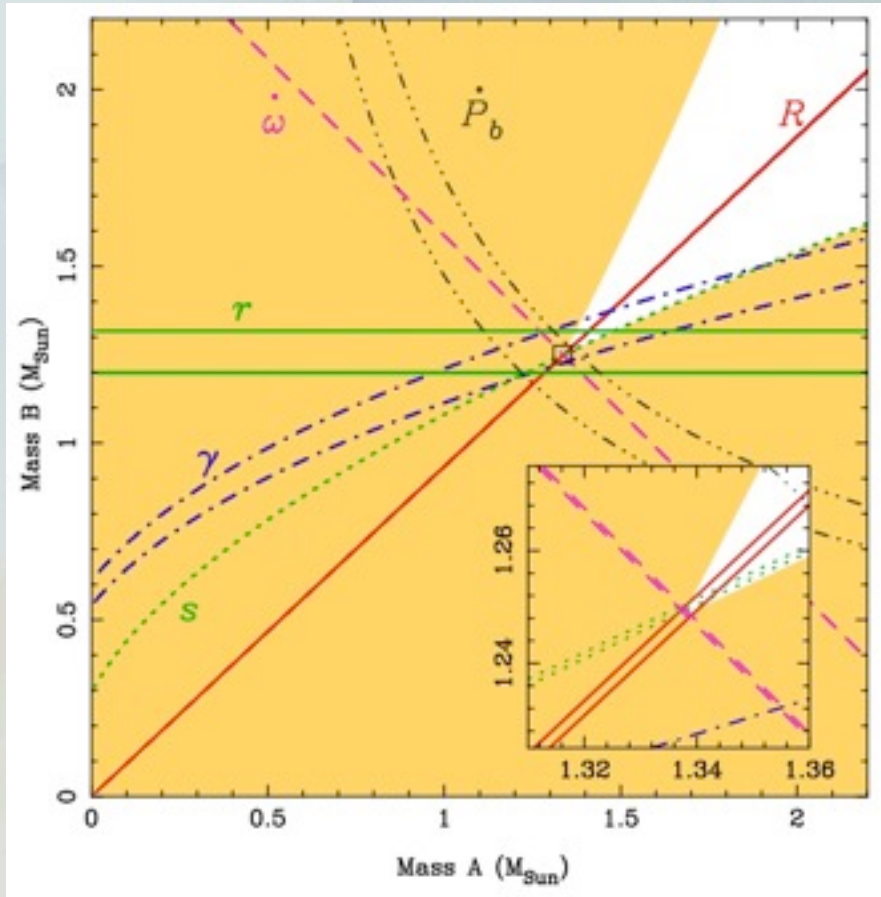
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## Orbital period decay

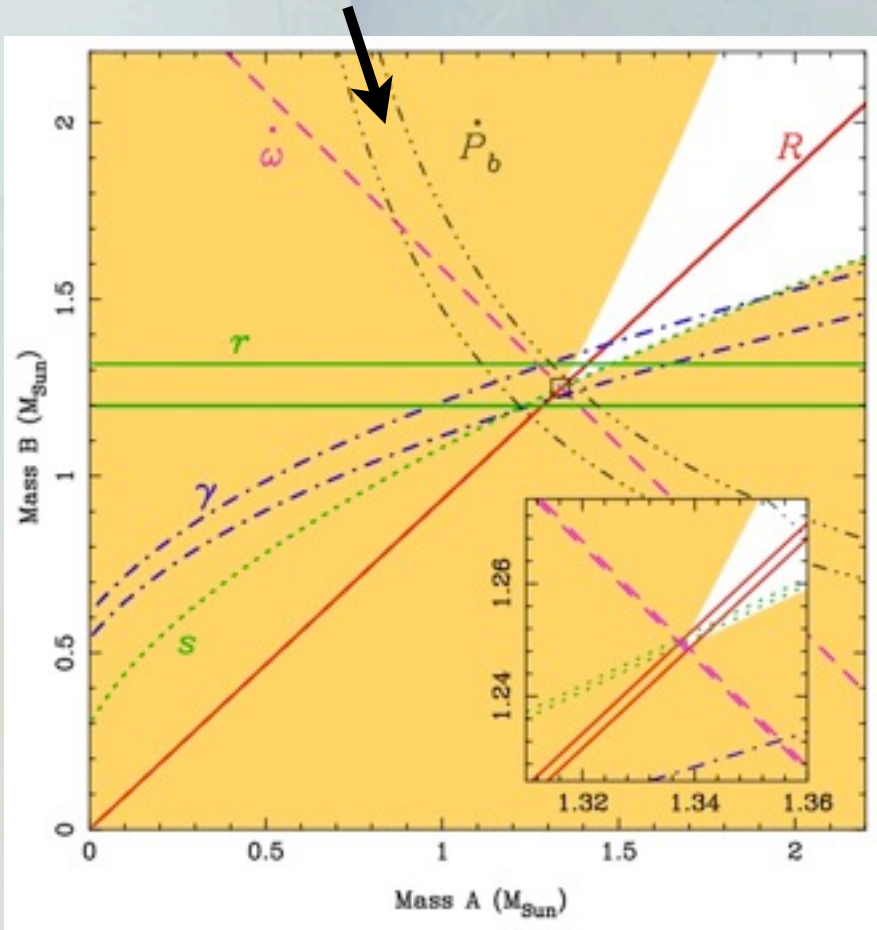
Gravitational wave emission

Orbit shrinks 7mm/day (coalescence timescale  $\sim 85$  Myrs)



# The Double Pulsar

## Testing gravity



Kramer et al., 2006

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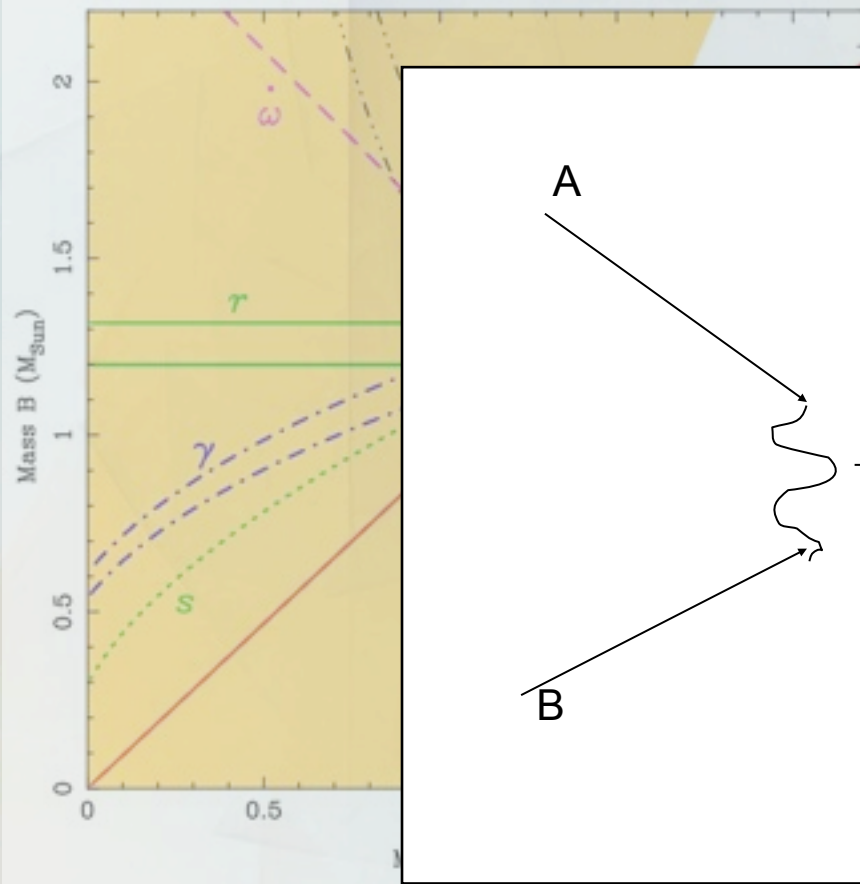
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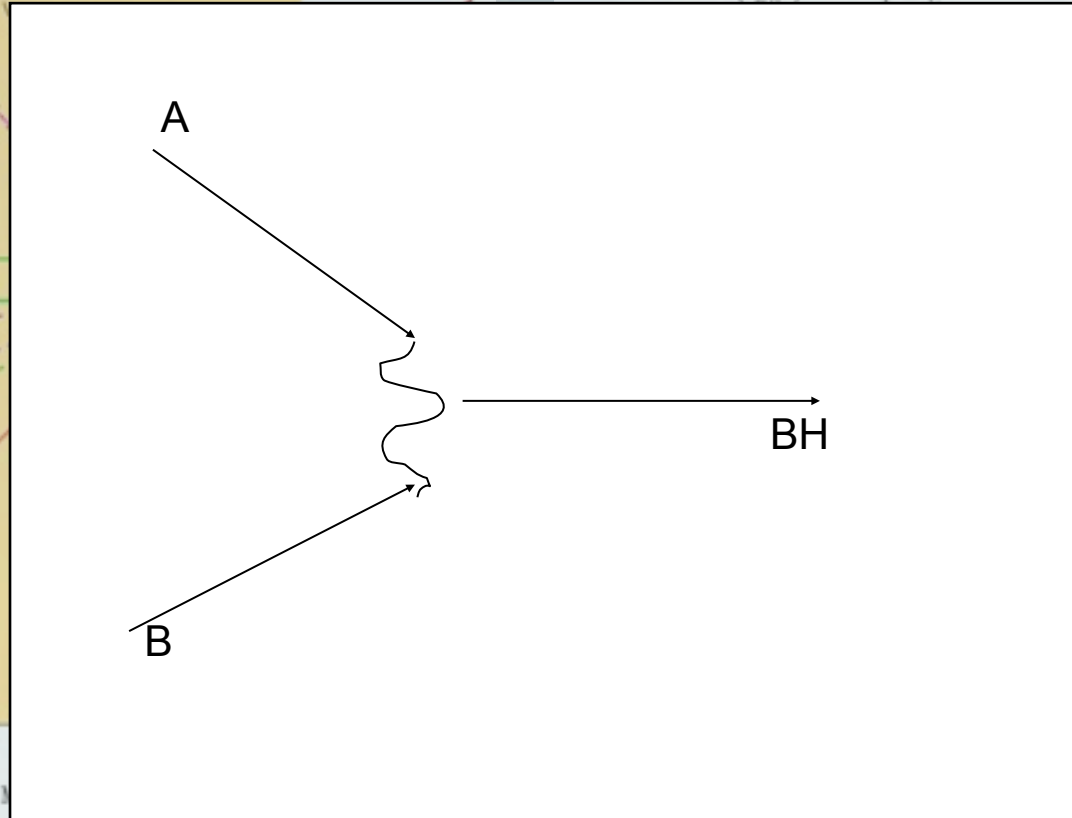
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$$\frac{37}{96} e^4 \frac{M_A M_B}{(M_A + M_B)^{1/3}},$$



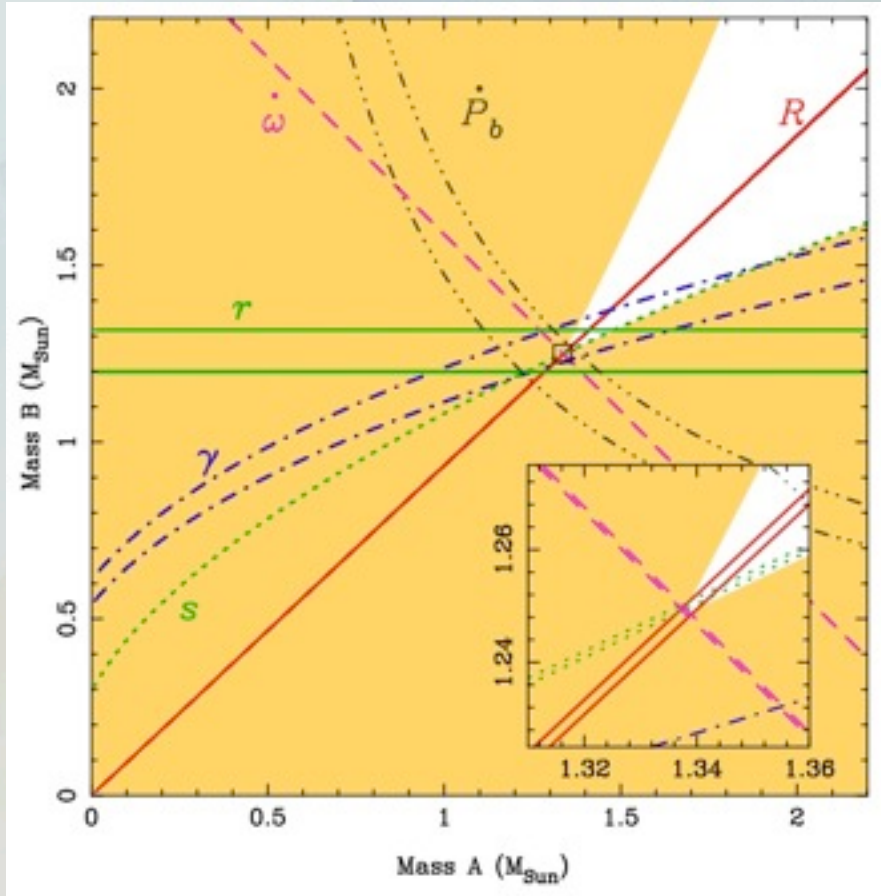
ssion

(coalescence

Kramer et al., 2006

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## Testing gravity



Kramer et al., 2006

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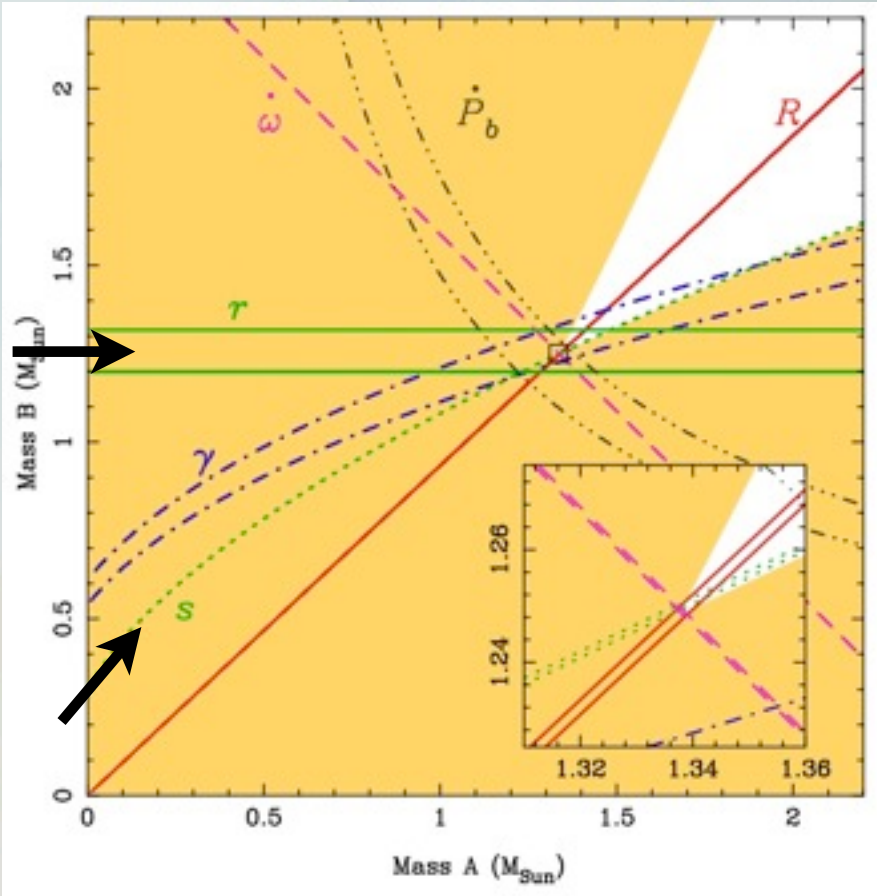
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## Shapiro delay

Light travels in curved space-time

# The Double Pulsar

## Testing gravity



Kramer et al., 2006

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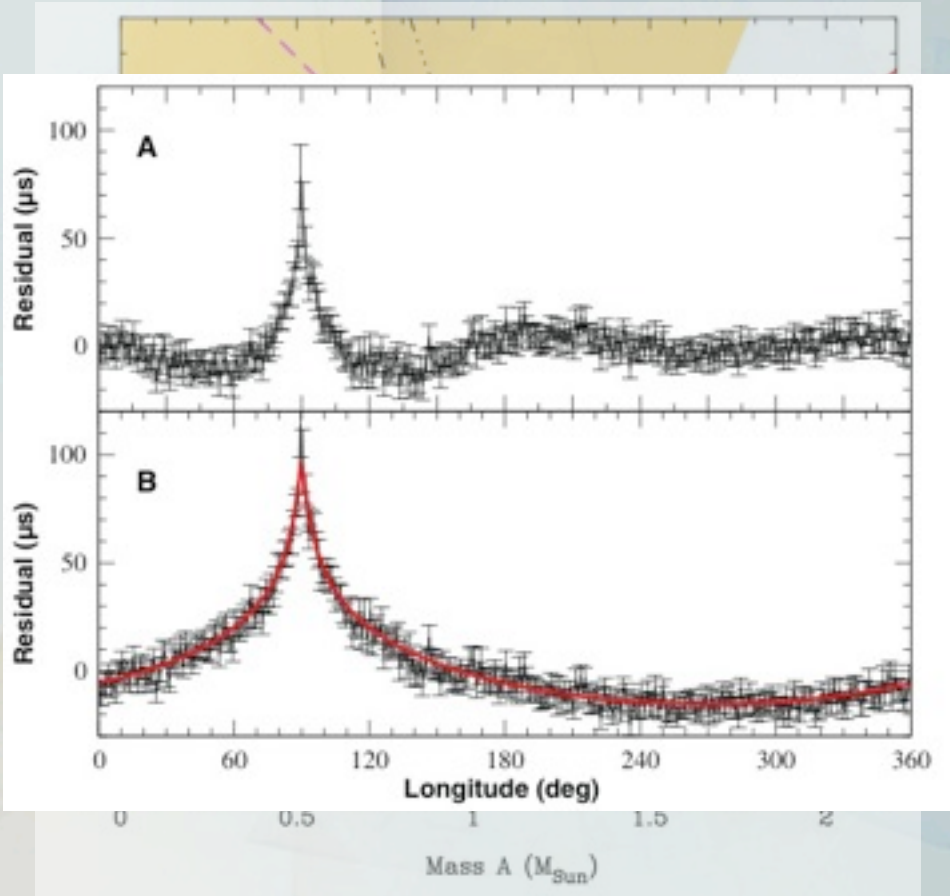
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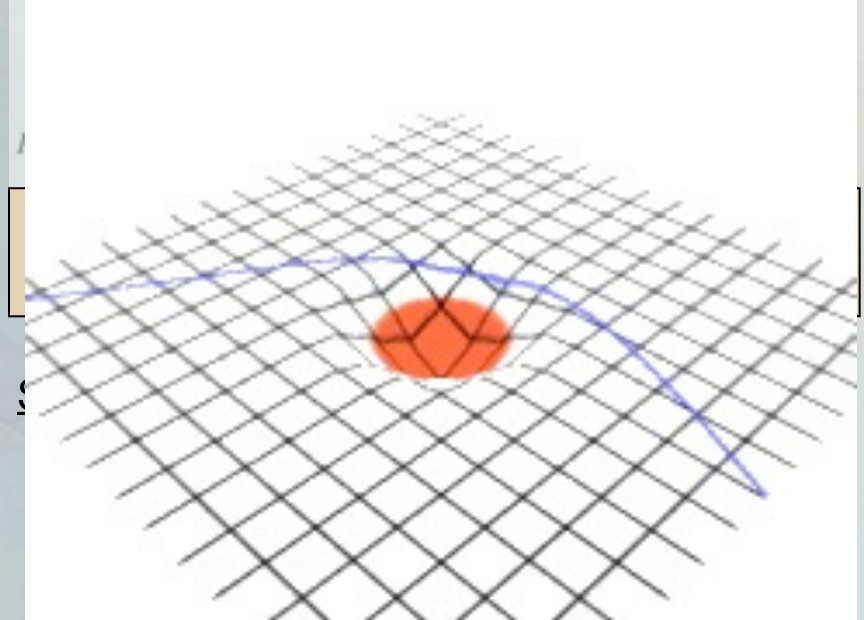
# The Double Pulsar

## Testing gravity



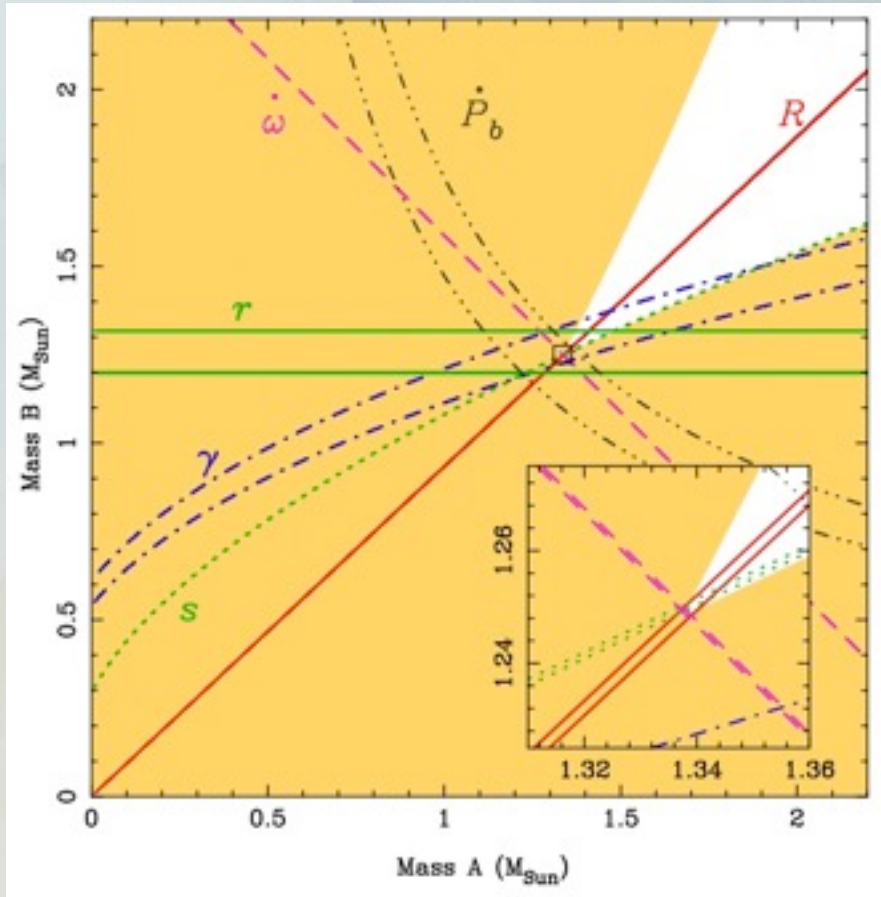
Kramer et al., 2006

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# The Double Pulsar

## Testing gravity



Kramer et al., 2006

## Mass ratio

*Double Pulsar only:*

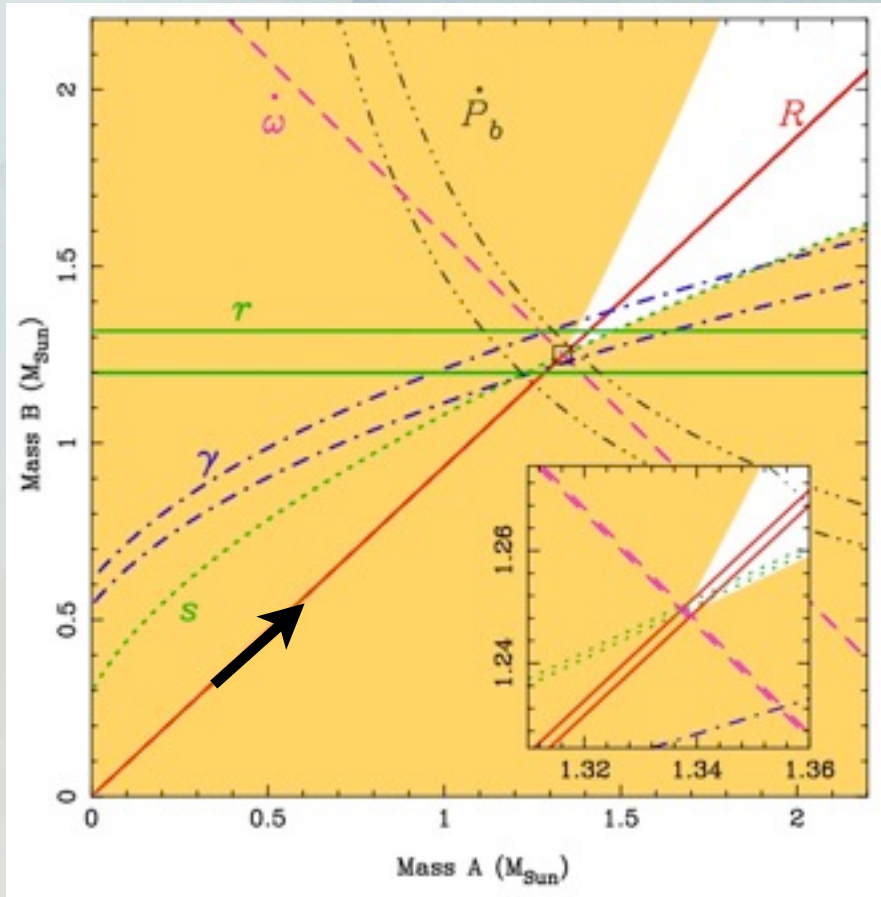
**mass ratio**,  $R$ .

$$M_A/M_B = a_B \sin(i) / a_A \sin(i)$$

(theory-independent; at 1PN level at least)

# The Double Pulsar

## Testing gravity



Kramer et al., 2006

## Mass ratio

*Double Pulsar only:*

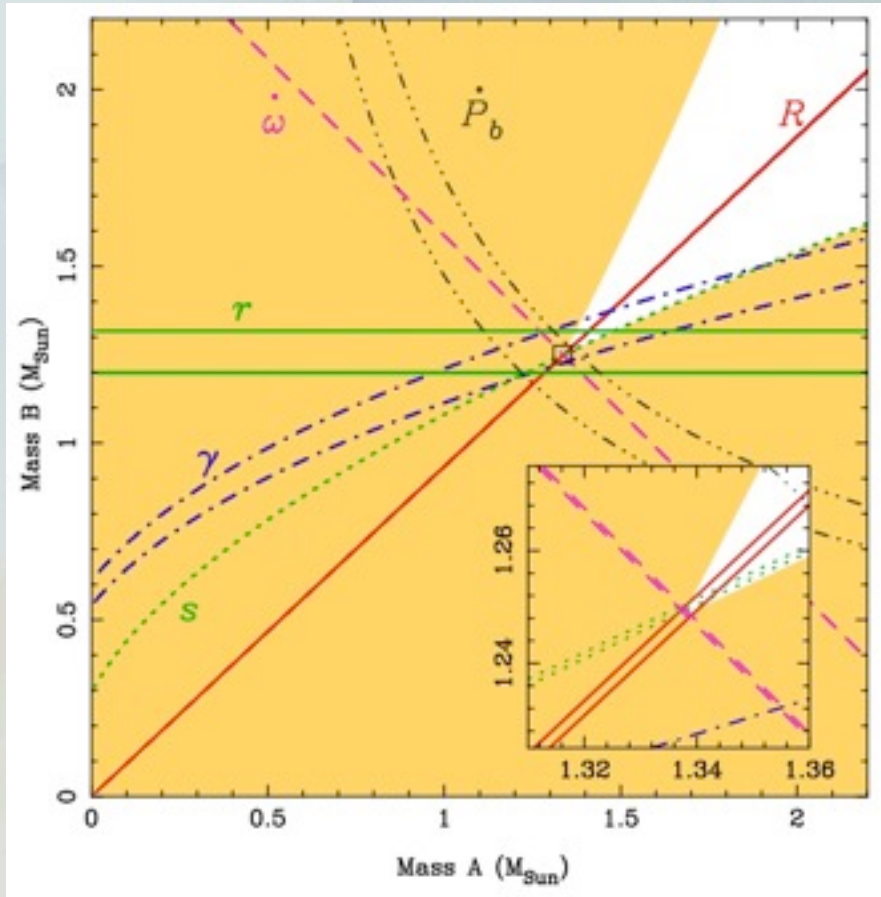
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# The Double Pulsar

## Testing gravity



Kramer et al., 2006

## Testing GR and gravity

Mass ratio + 5 PK parameters

=> Yields  $6-2=4$  gravity tests

## Best binary pulsar test

Shapiro delay “s” parameter  
Agrees with GR within 0.05%  
(Kramer et al. 2006)

Hulse-Taylor pulsar is 0.2%



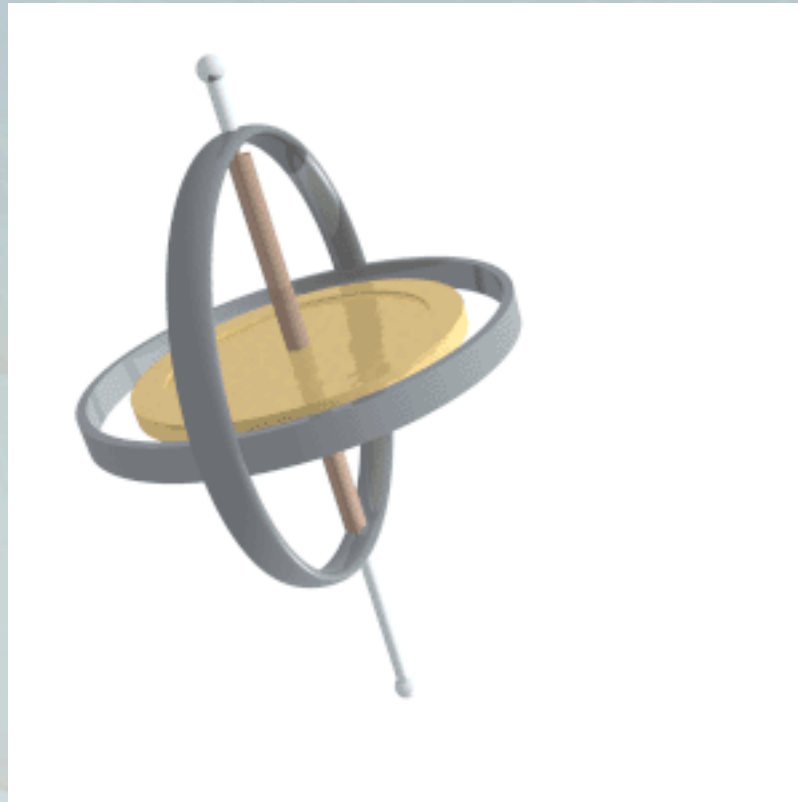
# Relativistic Spin Precession

Spin-orbit coupling induces relativistic spin precession

Precession of the spin angular momentum of a body is expected in relativistic systems.

Parallel transport in a curve space-time changes the orientation of a vector relative to distant observers.

Known as **“geodetic precession”** or “de Sitter – Fokker precession”.



# Relativistic Spin Precession

## Observational evidence of precession

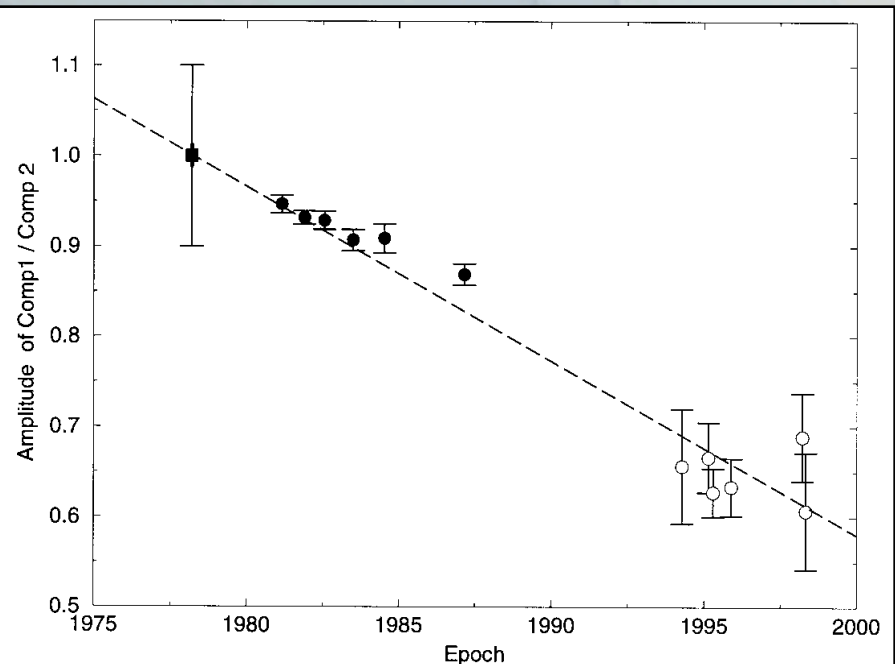
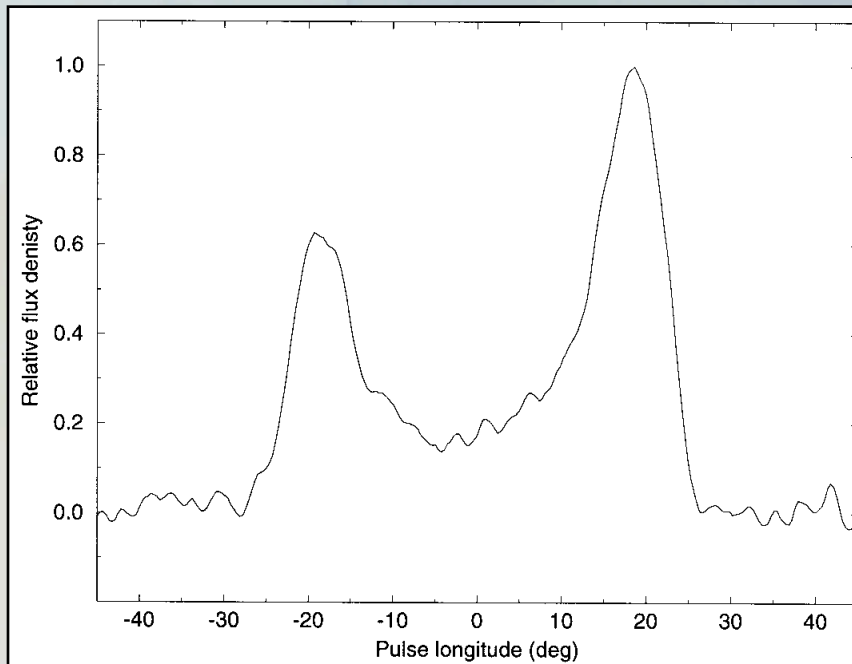
Changes of pulse shape and polarization over time.

PSR B1913+16 (Kramer 1998, Weisberg & Taylor 2002, Clifton & Weisberg 2008):

- Assume the predicted GR rate to map the beam

PSR B1534+12 (Stairs et al. 2004):

$$\bullet \Omega_{\text{obs}} = 0.44^{+0.48(4.6)}_{-0.16(0.24)} \text{ } ^\circ/\text{yr} \quad \Omega_{\text{GR}} = 0.51^\circ/\text{yr}$$



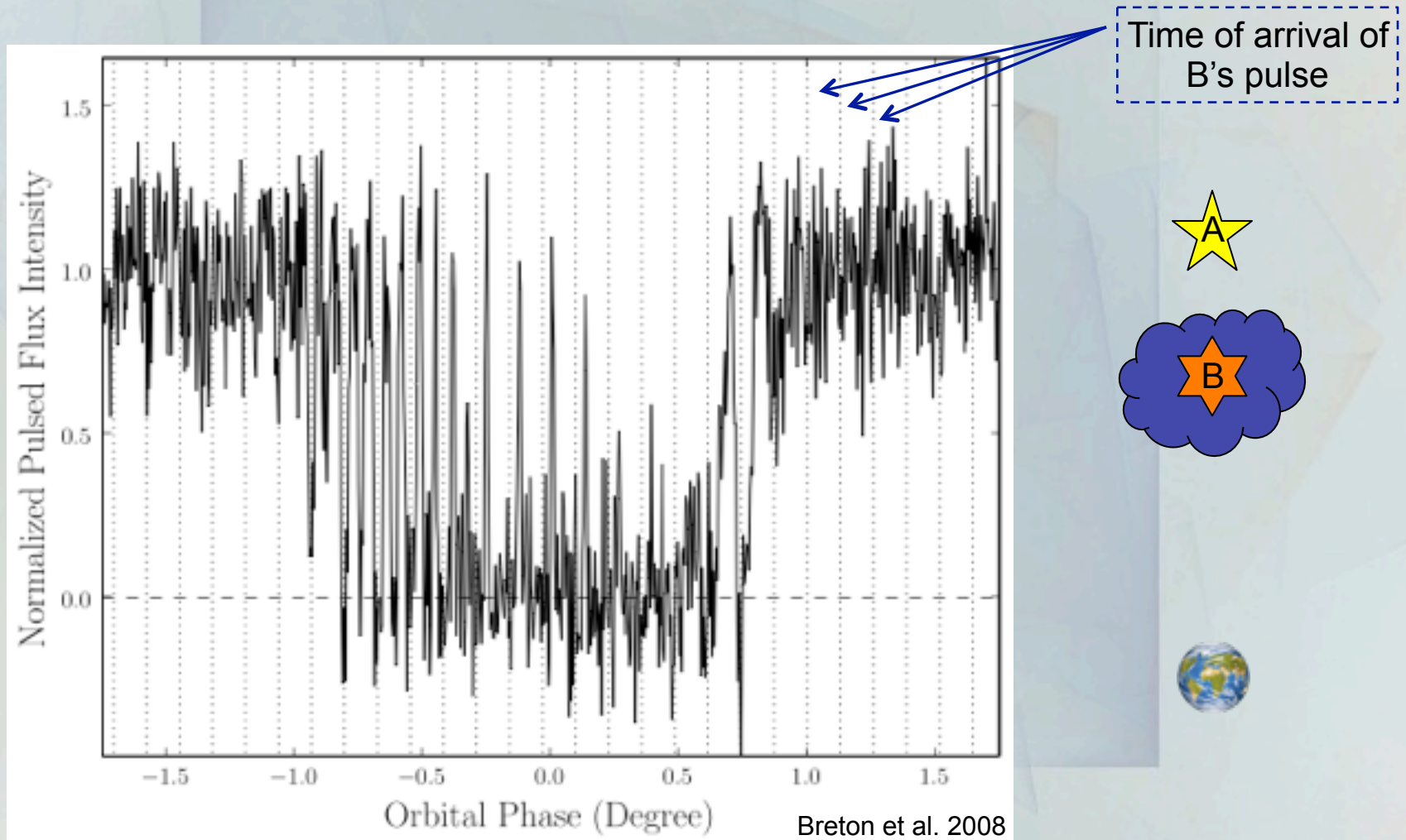
Kramer 1998

# Double Pulsar Eclipses

## Phenomenology

Pulsar A eclipsed by pulsar B for ~30 seconds at conjunction

**Rapid flux changes** synchronized with pulsar B's rotation (McLaughlin et al. 2004)



# Double Pulsar Eclipses

Modeling using synchrotron absorption in a truncated dipole

The “doughnut model” (Lyutikov & Thompson 2006)

# Double Pulsar Eclipses

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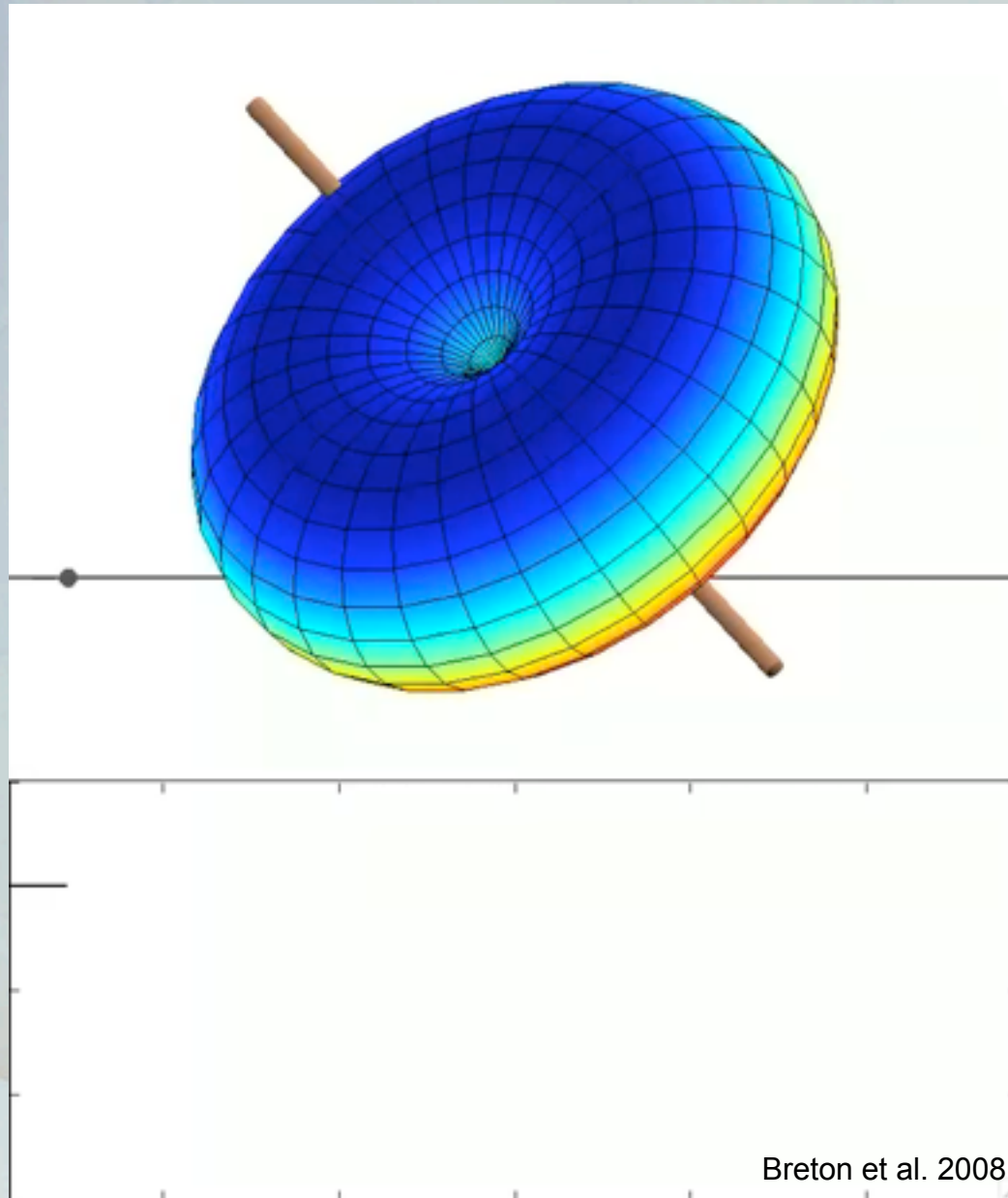
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Breton et al. 2008

# Double Pulsar Eclipses

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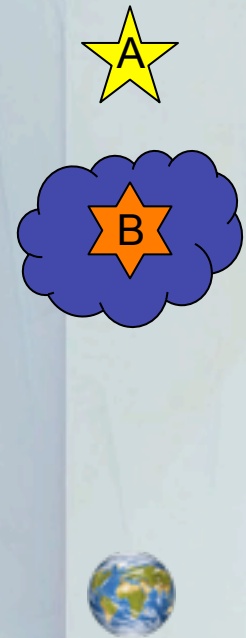
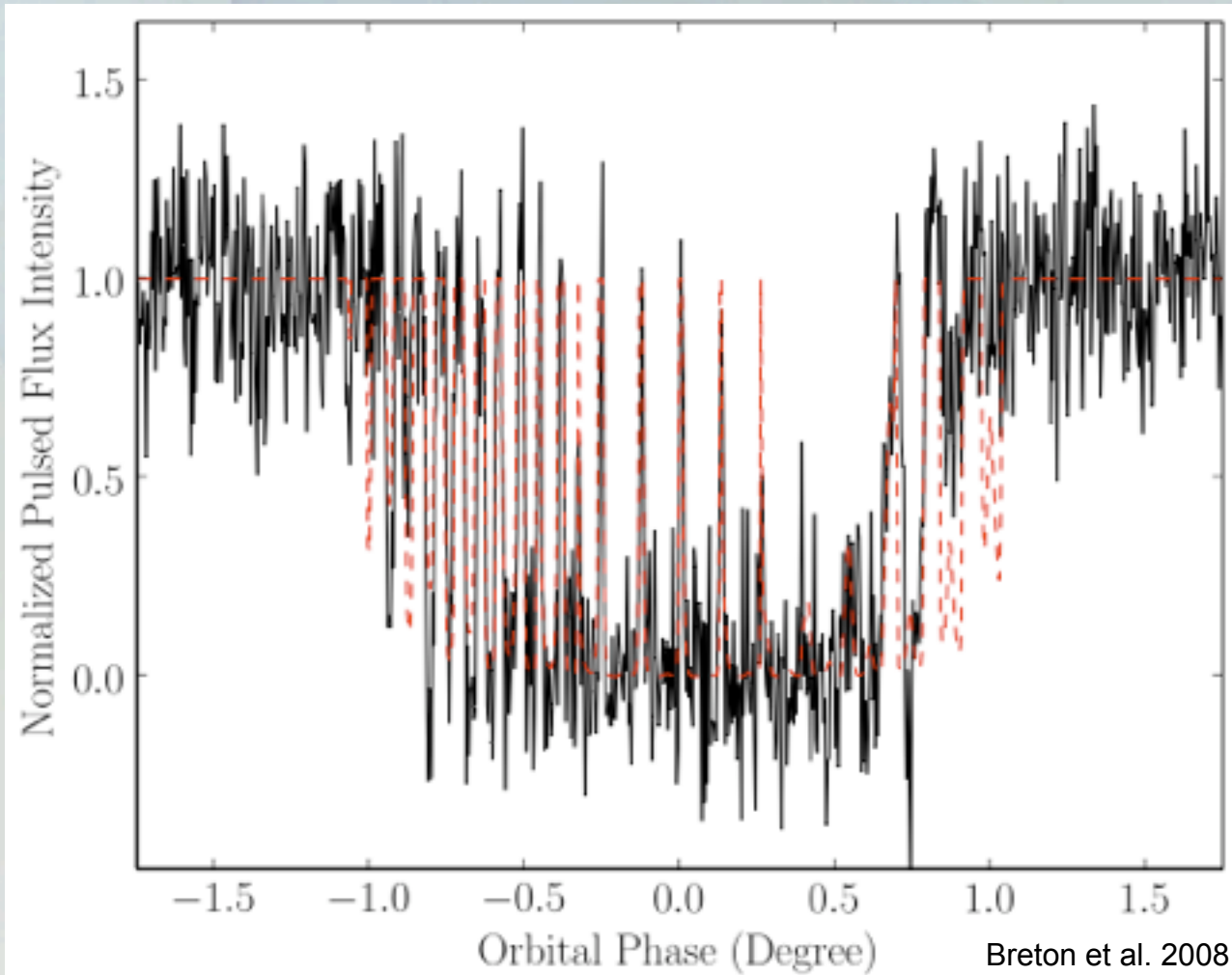


# Double Pulsar Eclipses

It works!

Empirical **evidence of dipolar magnetic** field geometry.

Eclipses are a proxy to infer the orientation of pulsar B.





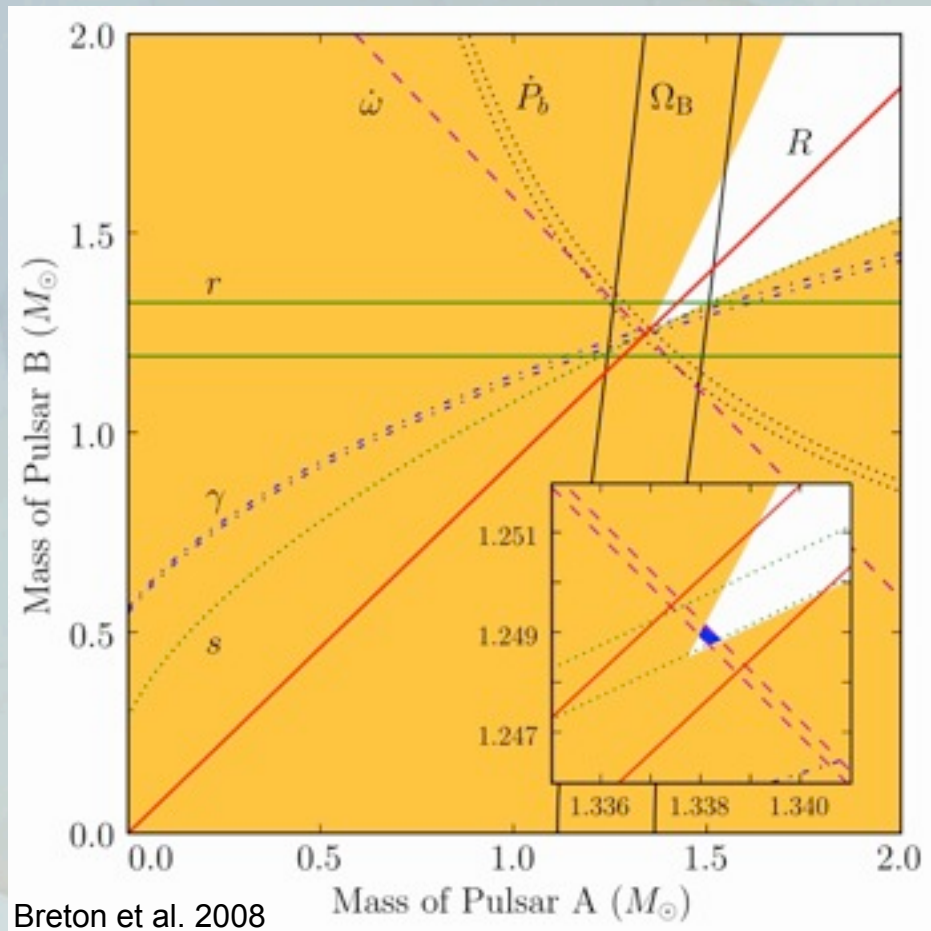
# Testing General Relativity in a Strong Field

General relativity passes the test once again!

Quantitative measurement of relativistic spin precession.

$$\Omega_B = \frac{1}{2} \left( \frac{2\pi}{P_{\text{orb}}} \right)^{5/3} \frac{m_A(4m_B + 3m_A)}{(1 - e^2)(m_B + m_A)^{4/3}} T_{\odot}^{2/3}$$

GR predicts:  $5^{\circ}07 \text{ yr}^{-1}$   
 Measurement:  $4^{\circ}77^{+0.66}_{-0.65} \text{ yr}^{-1}$



Breton et al. 2008

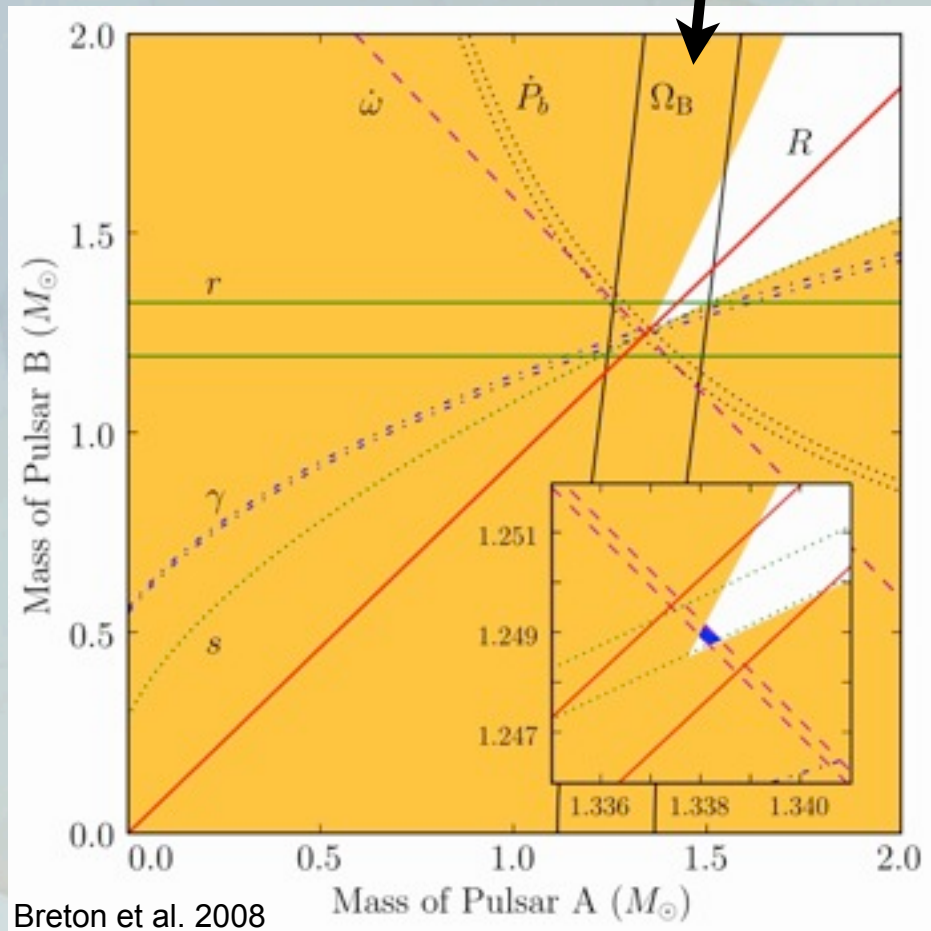
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By the way...

Pulsar B **disappeared in 2009**... so no more “Double Pulsar”

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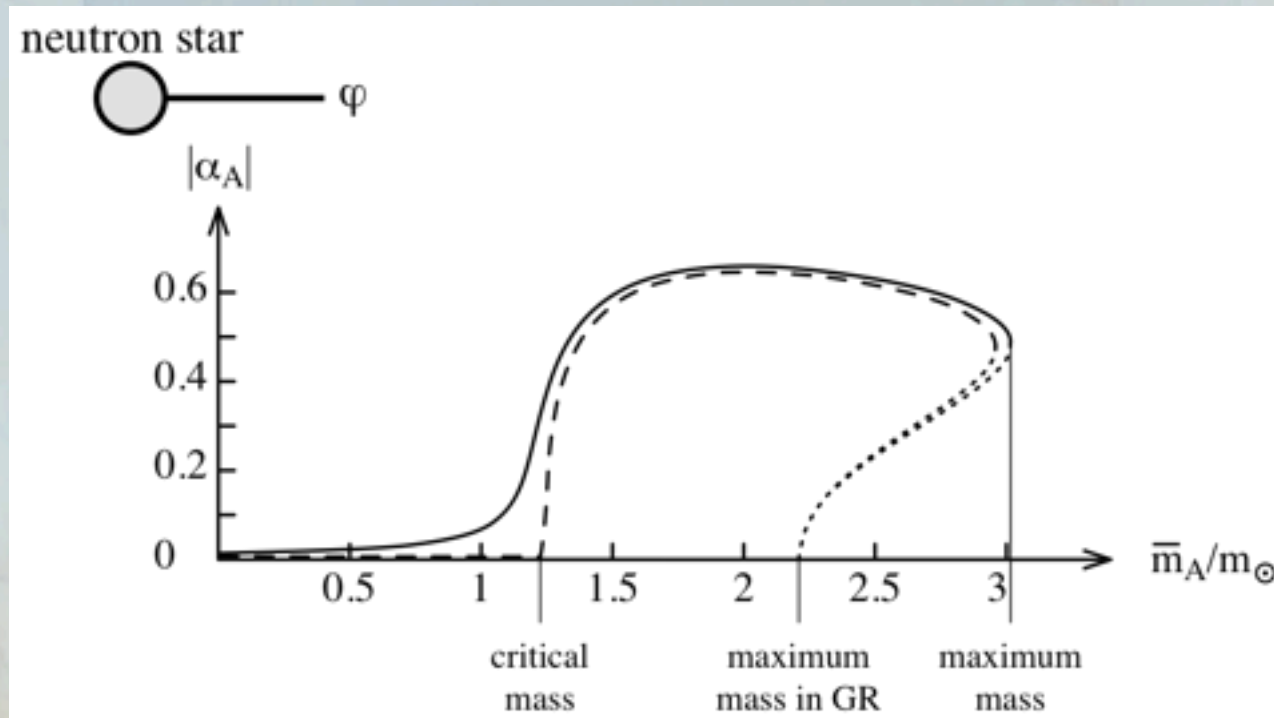
# Testing Theories of Gravity

# Strong-field regime

Why are tests of gravity involving pulsars interesting?

“Strong-field” regime:

- Still  $> 250\,000 R_{\text{Schwarzchild}}$  in separation but...
- $E_{\text{binding}} \sim 15\% E_{\text{tot}}$



Esposito-Farese (2004)

# Theory-Independent Test of Gravity

Unique test... for a unique double pulsar...

Generalized precession rate is (Damour & Taylor, 1992, Phys. Rev. D):

$$\Omega_B = \frac{\sigma_B L}{a_R^3 (1 - e^2)^{3/2}}$$

In terms of a very particular choice of observable timing parameters:

$$\Omega_B = \frac{x_A x_B}{s^2} \frac{n^3}{1 - e^2} \frac{c^2 \sigma_B}{\mathcal{G}}$$

# Theory-Independent Test of Gravity

Unique test... for a unique double pulsar...

Generalized precession rate is (Damour & Taylor, 1992, Phys. Rev. D):

$$\Omega_B = \frac{\sigma_B L}{a_R^3 (1 - e^2)^{3/2}}$$

In terms of a very particular choice of observable timing parameters:

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**Test of the strong-field parameters**  $\frac{\sigma_B}{\mathcal{G}}$ .

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The *Double Pulsar* is the only double neutron star system in which both semi-major axes can be **measured independently!**

# Dipolar Gravitational Waves and MOND

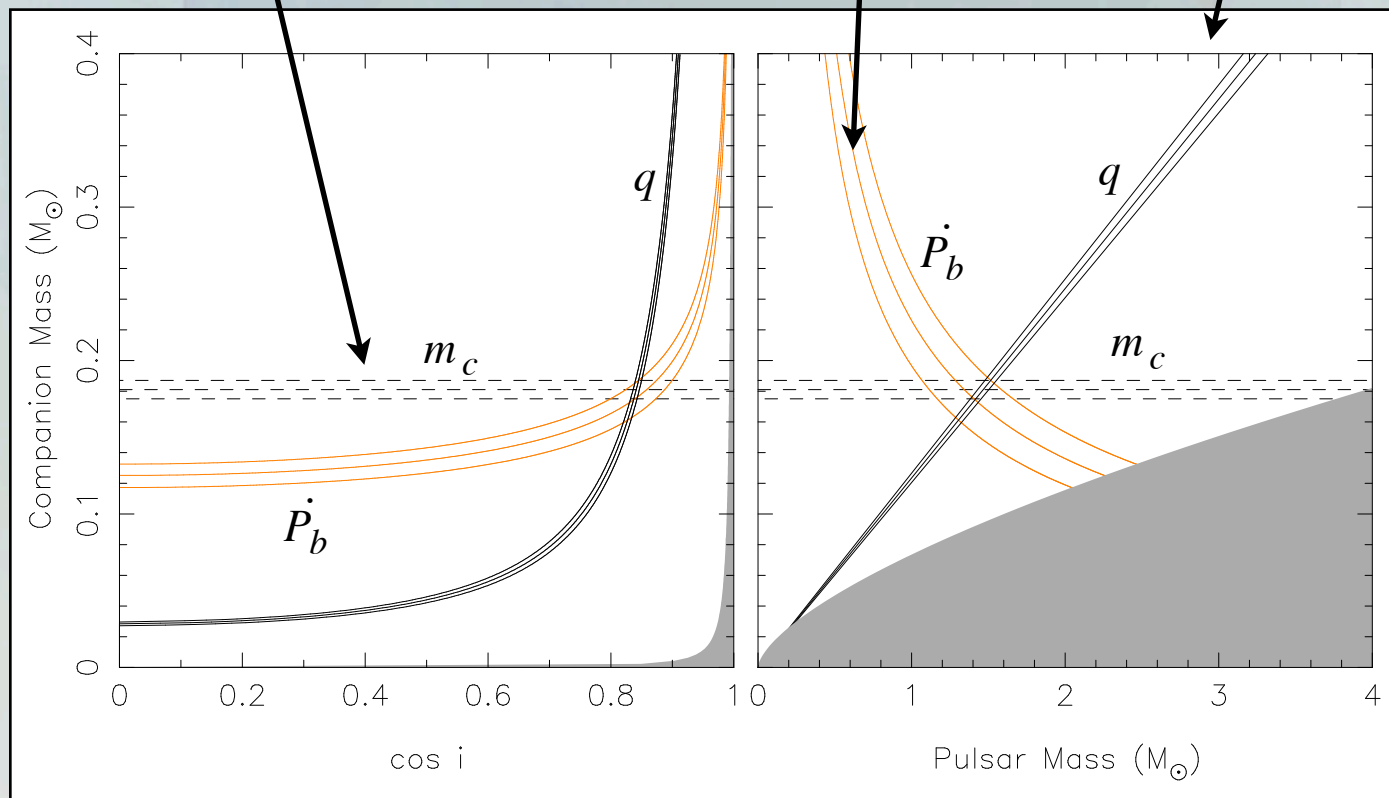
## The pulsar-white dwarf binary PSR J1738+0333

See Freire et al. (2012) & Antoniadis et al. (2012)

White dwarf mass from photometry + spectroscopy

Mass ratio from pulsar timing + white dwarf spectroscopy

Orbital period decay from pulsar timing



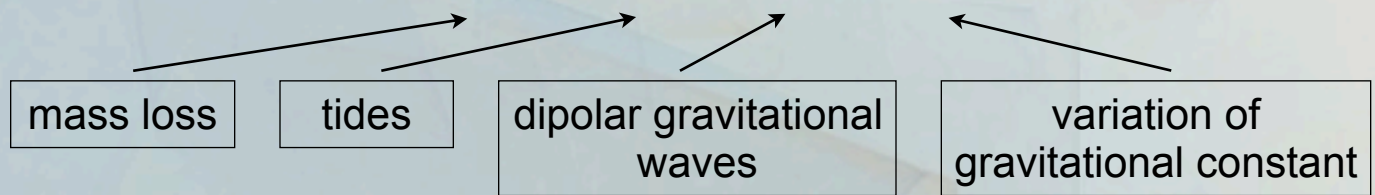
Freire et al. (2012) & Antoniadis et al. (2012)

# Dipolar Gravitational Waves and MOND

## Constraining dipolar gravitational waves

The “excess” orbital period decay:

$$\dot{P}_b^{\text{XS}} = \dot{P}_b^{\dot{M}} + \dot{P}_b^{\text{T}} + \dot{P}_b^{\text{D}} + \dot{P}_b^{\dot{G}}$$



$$\dot{P}_b^{\text{D}} \simeq -2\pi n_b \mathbb{T}_{\odot} m_c \frac{q}{q+1} \kappa_D (s_{\text{psr}} - s_{\text{comp}})^2$$

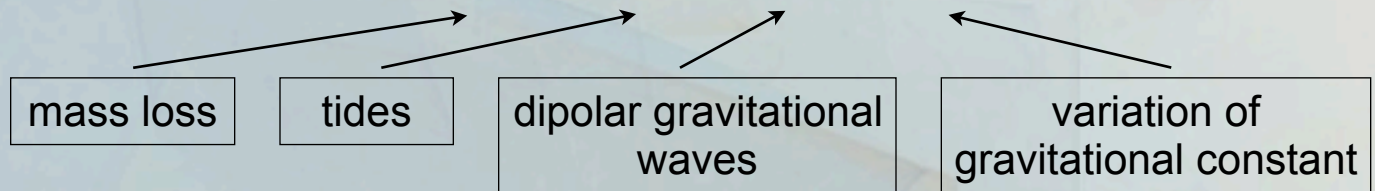
where  $s_{\text{psr,comp}} \simeq \epsilon / M_{\text{psr,comp}} c^2$

# Dipolar Gravitational Waves and MOND

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Neutron star + white dwarf binaries  
have very **different self-gravity**

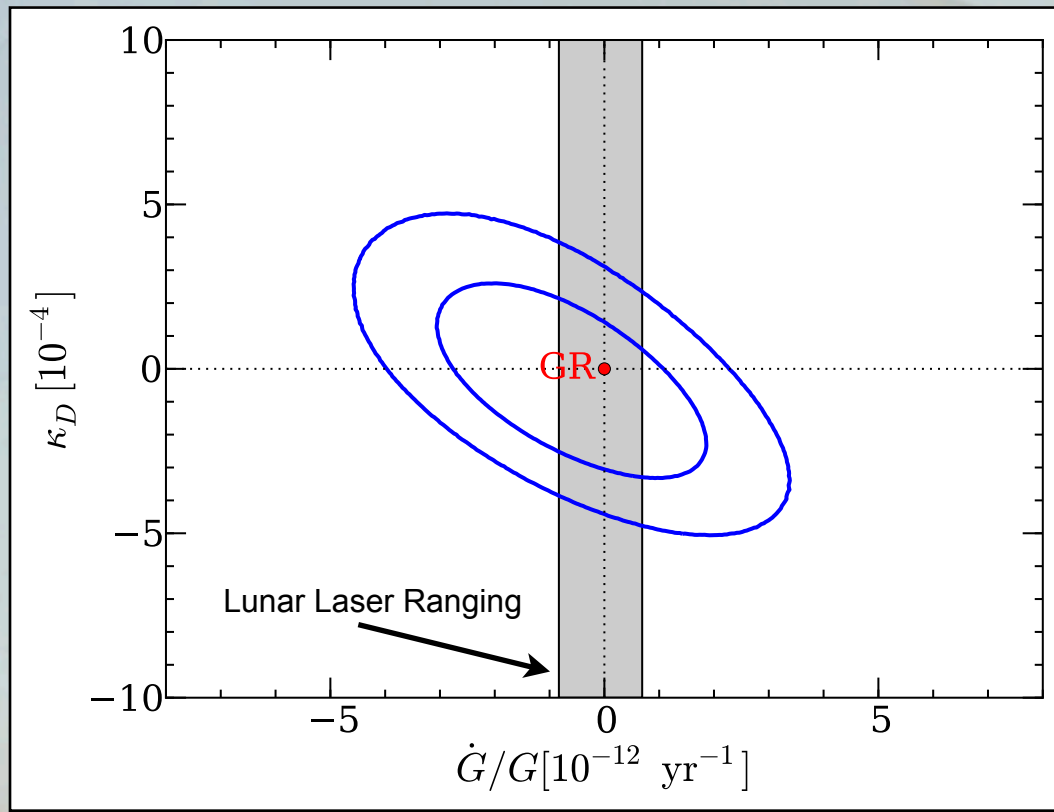
$$s_{\text{psr}} \neq s_{\text{comp}}$$

# Dipolar Gravitational Waves and MOND

## Constraining dipolar gravitational waves

$\dot{P}_b^D$  and  $\dot{P}_b^{\dot{G}}$  are degenerate

Combining data from PSRs J1738+0333 and J0437-4715 breaks the degeneracy



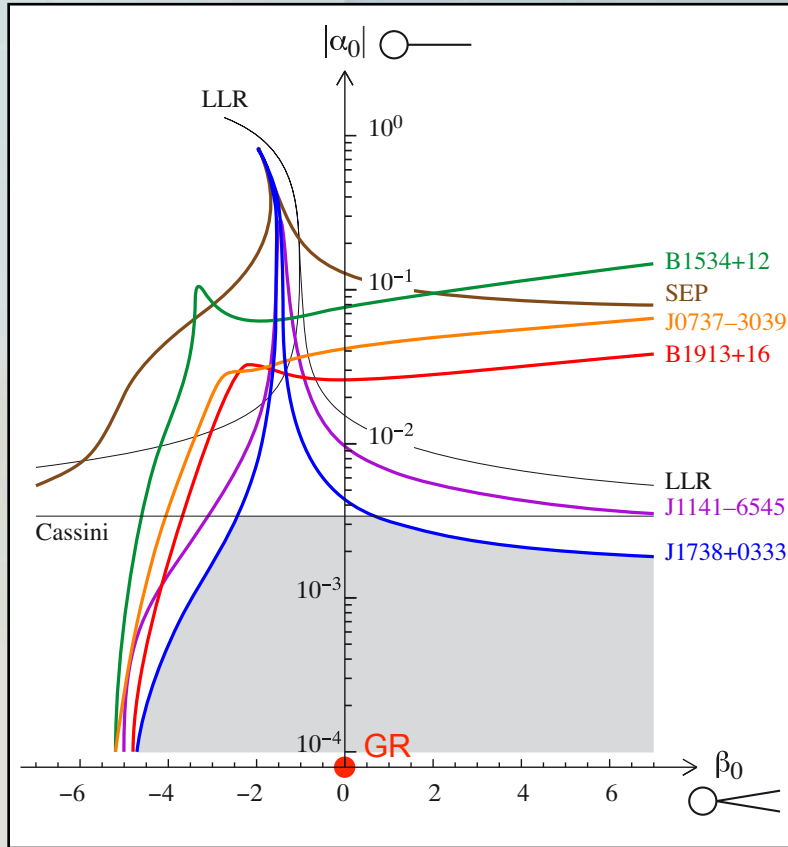
Freire et al. (2012) and, also, Lazaridis et al. (2009)

# Dipolar Gravitational Waves and MOND

## Constraining tensor-scalar theories

A general tensor scalar field  $\ln A(\varphi) - \ln A(\varphi_0) = \alpha_0(\varphi - \varphi_0) + \frac{1}{2}\beta_0(\varphi - \varphi_0)^2$

Associated metric  $\tilde{g}_{\mu\nu} \equiv A^2(\varphi)g_{\mu\nu}$



Freire et al. (2012)

Modify the behavior of gravitational constant

$$\tilde{G}_{AB} \equiv G_* A^2(\varphi_0) \cdot (1 + \alpha_A \alpha_B)$$

And the Eddington's PPN parameters

$$\gamma_{AB} \equiv 1 - 2 \frac{\alpha_A \alpha_B}{1 + \alpha_A \alpha_B}$$

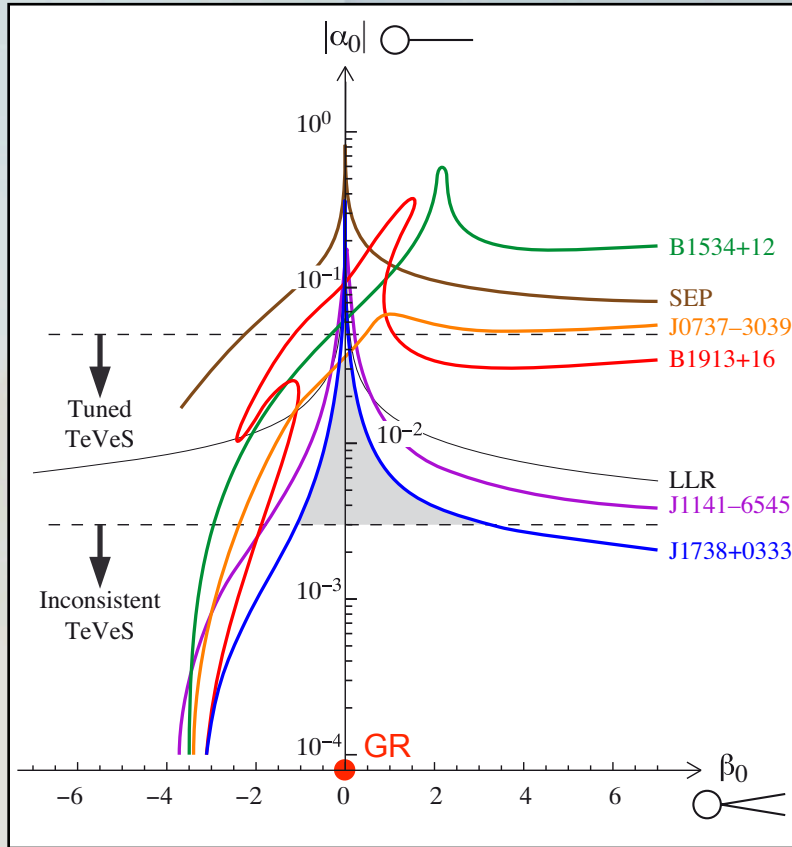
$$\beta_{BC}^A \equiv 1 + \frac{1}{2} \frac{\beta_A \alpha_B \alpha_C}{(1 + \alpha_A \alpha_B)(1 + \alpha_A \alpha_C)}$$

# Dipolar Gravitational Waves and MOND

## Constraining modified newtonian dynamics

MOND is the non-relativistic limit of tensor-vector-scalar (TeVeS) theories

Associated metric  $\tilde{g}_{00} = A^2(\varphi)g_{00}$  and  $\tilde{g}_{ij} = A^{-2}(\varphi)g_{ij}$



Freire et al. (2012)

Modify the behavior of gravitational constant

$$\tilde{G}_{AB} = G_*(1 + \alpha_A \alpha_B)$$

And the Eddington's PPN parameters

$$\gamma_{AB} = \gamma^{\text{PPN}} = 1$$

$$\beta_{BC}^A \equiv 1 + \frac{1}{2} \frac{\beta_A \alpha_B \alpha_C}{(1 + \alpha_A \alpha_B)(1 + \alpha_A \alpha_C)}$$

# More Tests of Gravity?

## SEP violation, preferred-frame effects

Pulsar + WD in low eccentricity wide orbit

Test strong equivalence principle  $P_{\text{orb}}^2/e^2$   $|\Delta| < 4.6 \times 10^{-3}$

Test violation of momentum conservation and preferred frame-effects

$$P_{\text{orb}}^2/(eP) \quad P_{\text{orb}}^{1/3}/e$$

statistical: see, e.g., Gonzalez et al. 2011, Stairs et al. 2005

$$|\hat{\alpha}_3| < 5.5 \times 10^{-20}$$

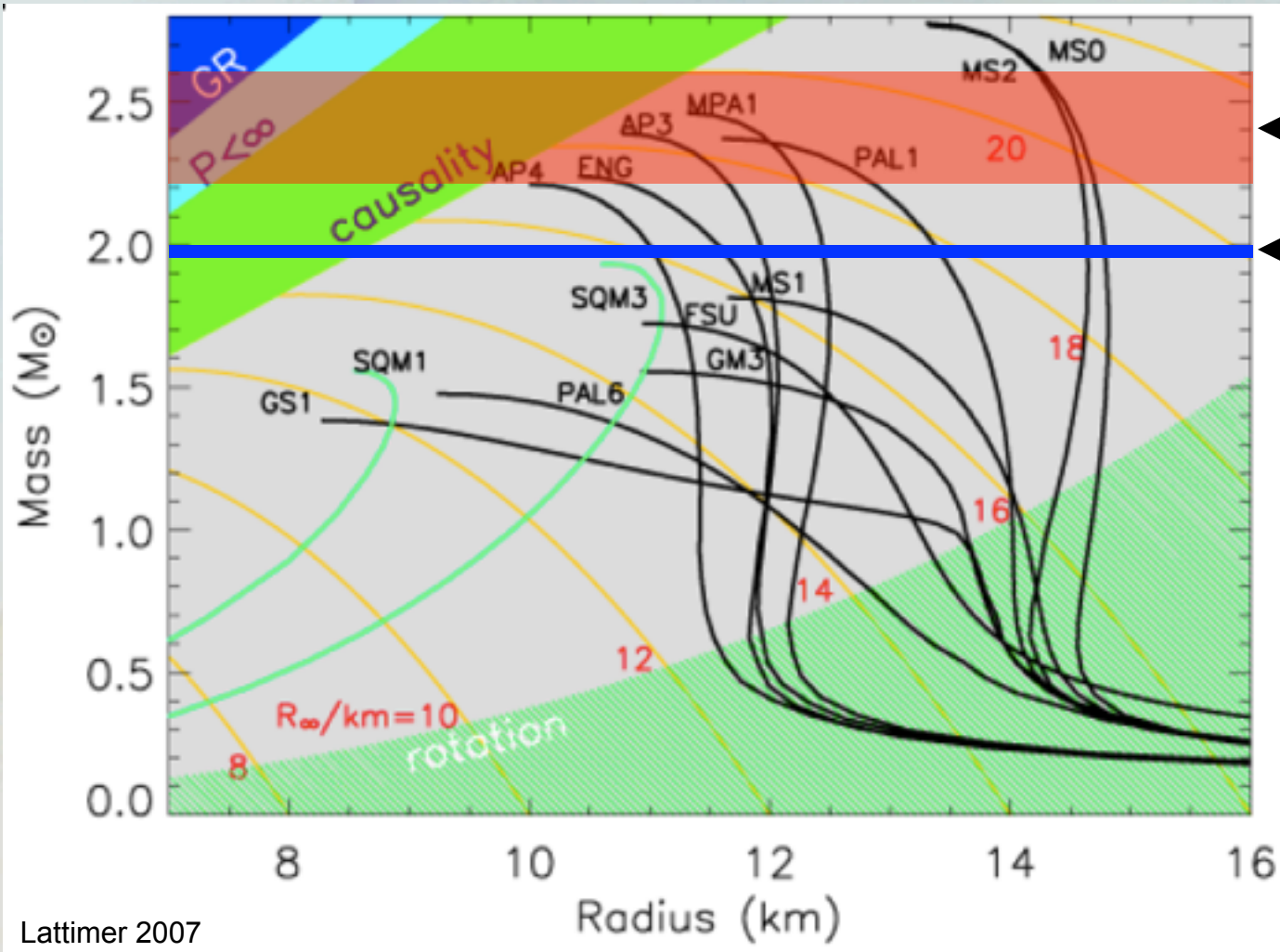
individual system: see, e.g., Wex & Kramer 2007





# Neutron Star Equation-of-State

Probing the upper mass limit of neutron stars



van Kerkwijk, Breton  
& Kulkarni (2011)

← PSR B1957+20

← PSR J1614-2230

Demorest et al. (2010)

Lattimer 2007

# The Future Is Bright

More pulsars, more photons

More pulsars: What if we know “all” pulsars from the Galaxy?

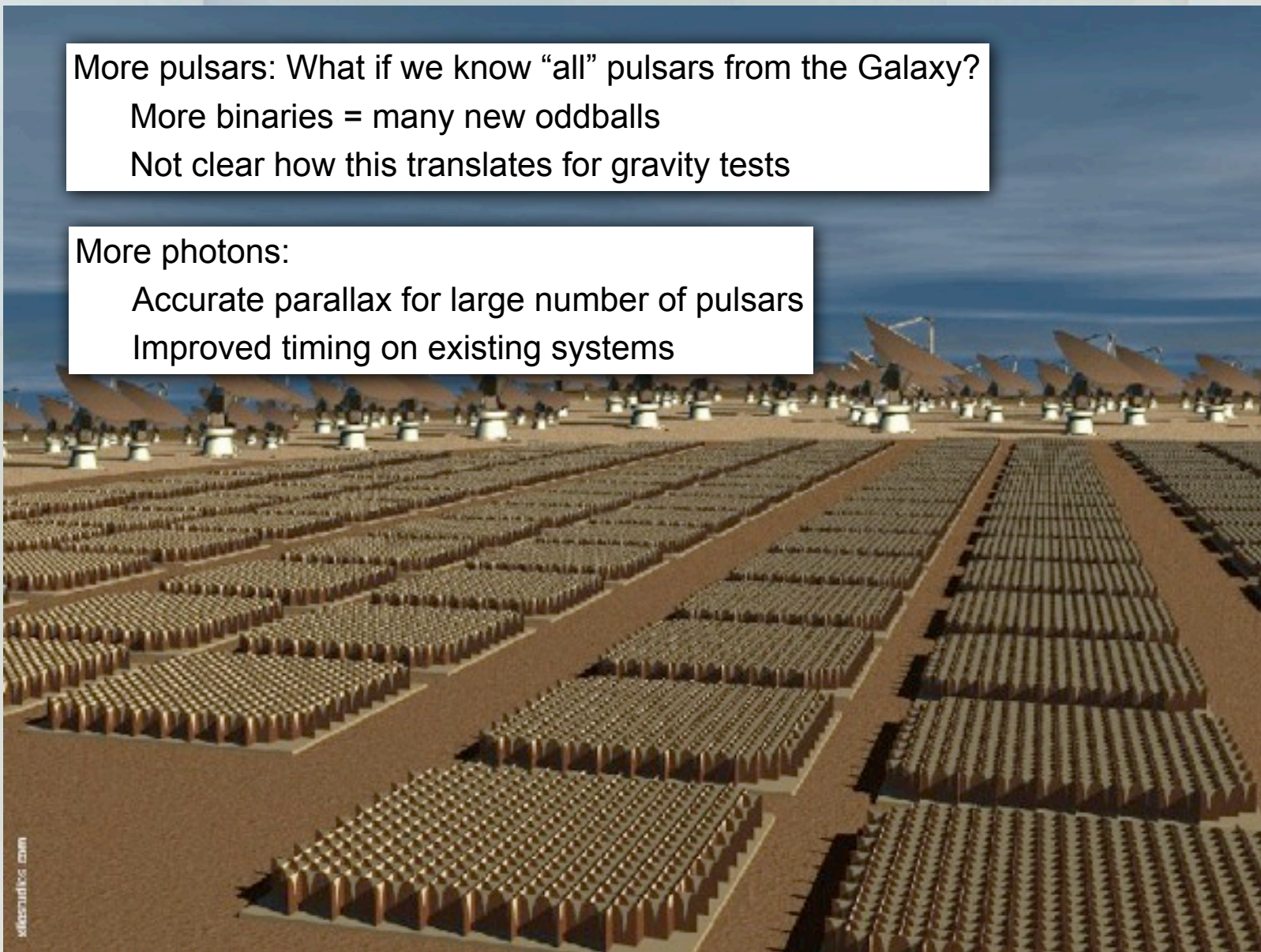
More binaries = many new oddballs

Not clear how this translates for gravity tests

More photons:

Accurate parallax for large number of pulsars

Improved timing on existing systems



skasouthafrica.com

# Conclusions

